



Physical Mechanisms at Play In the Multi-Scales Interaction Between MHD and Turbulence

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In collaboration with

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Magnetic islands are ubiquitous in nature

On Solar Surface

[September 2005, captured in the X-ray waveban by NASA's TRACE satellite. Photo courtesy of the University of California Berkeley, all rights reserved]



In tokamak



Current driven instability



Magnetic islands degrade tokamak confinement

Growth of Neoclassical Tearing Modes (NTMs) : magnetic island degrading the plasma pressure and sometimes causing disruption

Behavior of NTMs : modified Rutherford equation



[R.J. La Haye, POP 13 (2006)]

NTMs precursors :

♦ Sawtooth oscillations
♦ Fishbones instabilities
♦ Edge localized modes
♦ ... ?

 In JT-60U, 80% of high β discharges, a (2/1) NTM appears without precursor event [A. Isayama et al, PFR 8 (2013)]

Open question : origin of seed island ?

MHD-Turbulence Interaction, a Multi-Scales Problem



 Interchange like instabilities coexist with macro MHD instabilities and lead to micro-turbulence in fusion devices.

The interaction of magnetic island with interchange is a multi-scales problem.

[F. Militello et al, POP 15 (2008)]
[F.L. Waelbroeck et al, PPCF 51 (2009)]
[M. Muraglia et al, PRL 103 (2009)]
[A. Ishizawa et al, POP 17 (2010)]
[F. Hariri et al, PPCF 57 (2015)]
[L. Bardoczi et al, POP 24 (2017)]

Turbulence Driven Magnetic Island (TDMI)

[M. Muraglia et al, PRL 107 (2011)]
[A. Poyé et al, POP 22 (2015)]
[W. Hornsby et al, PPCF 58 (2015)]
[O. Agullo et al, POP 24 (2017)]
[A. Ishizawa et al, PPCF 61 (2019)]

TDMI amplified and at the origin of a NTM

[M. Muraglia et al, NF (2017)] J. Frank, PhD thesis, CEA and PIIM Lab

Outline

I. Introduction MHD-Turbulence Interaction Modelisation

II. Turbulence driven island in a 2D slab geometry (M. Muraglia)

III. Nonlinear growth of NTM from a seed turbulence driven island

(M. Muraglia)

IV. Coherent and non local beating leading to island generation in a 3D geometry (N. Dubuit)

V. Conclusions

I. Magnetic island : Tearing Instability



I. Interchange Instability



Analogy : Rayleigh-Taylor instability



<u>Tokamak</u> :

- Curvature and pressure gradient in opposition => Interchange instability
- => Turbulence



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I. Model: Reduced MHD

Minimal Model :

- 2D Slab => near a resonant surface in a (x, y) poloidal cross section (mono-helicity)
- Reduced MHD => Fluctuations dynamic evolution of electrostatic potential ϕ , electronic pressure *p* and magnetic flux ψ .
- Minimal Interchange Model :

Momentum conservation

$$\partial_t \nabla^2_\perp \phi + [\phi, \nabla^2_\perp \phi] = -\kappa_1 \partial_y P + \nu \nabla^4_\perp \phi$$

Energy conservation

 $\partial_t P + [\phi, P] = -\partial_x P_0 \partial_y \phi + \chi_\perp \nabla_\perp^2 P$

Minimal Tearing Mode Model :
 Momentum conservation
 $\partial_t \nabla^2_{\perp} \phi + [\phi, \nabla^2_{\perp} \phi] = [\psi + \psi_0, \nabla^2_{\perp} \psi] + \nu \nabla^4_{\perp} \phi$

Ohm's law

$$\partial_t \psi = [\psi + \psi_0, \phi] + \eta \nabla_\perp^2 \psi$$

Model includes both resistive Interchange and Tearing Mode :

$$\partial_t \nabla_{\perp}^2 \phi + [\phi, \nabla_{\perp}^2 \phi] = [\psi + \psi_0, \nabla_{\perp}^2 \psi] - \kappa_1 \partial_y P + \nu \nabla_{\perp}^4 \phi$$

$$\partial_t P + [\phi, P] = -\partial_x P_0((1 - \kappa_2)\partial_y \phi + \kappa_2 \partial_y P) + \rho_{\star}^2 [\psi + \psi_0, \nabla_{\perp}^2 \psi] + \chi_{\perp} \nabla_{\perp}^2 P$$

$$\partial_t \psi = [\psi + \psi_0, \phi - P] - \partial_x P_0 \partial_y \psi + \eta \nabla_{\perp}^2 \psi$$

[M. Muraglia et al, NF 49, 055016 (2009)]

I. Model: Reduced MHD

- Numerical resolution of the model => Code AMON
 - temporal evolution : Runge-Kutta order 4th
 - semi-spectral code in a numerical box [Lx, Ly] finite difference in the radial x direction
 Fourrier decomposition in the poloidal y direction

$$\psi(x, y, t) = \sum_{m \in \mathbb{Z}} \psi_m(x, t) \exp^{i\frac{2\pi m}{Ly}y}$$

• Instability characterization : THE PARITY of the eigenfunctions ψ_m , ϕ_m , P_m



0.015



II. NL generation of TDMI

Linear Spectrum



 Δ^{\prime} Δ^{\prime} Δ^{\prime} Δ^{\prime} Linear spectrum is stable with respect with tearing instability

=> No island

- Stable large scales modes
- Small scales turbulence driven by interchange instability

=> What 's about non-linear dynamics ?

II. NL generation of TDMI



NL generation of TDMI by a beating of interchange modes

[M. Muraglia et al, PRL 107 (2011)] & [A. Poyé et al, POP 22 (2015)] [W. Hornsby et al, PPCF 58 (2015)]

II. Origine of the island



II. Origine of the island



III. NL growth of NTM ?



Is such turbulence driven seed island at the origin of NTM growth ?

=> New model including bootstrap current is required

In 2D slab geometry :

 $\partial_t \nabla_{\perp}^2 \phi + [\phi, \nabla_{\perp}^2 \phi] = [\psi + \psi_0, \nabla_{\perp}^2 \psi] - \kappa_1 \partial_y P + \nu \nabla_{\perp}^4 \phi$ $\partial_t P + [\phi, P] = -\partial_x P_0((1 - \kappa_2)\partial_y \phi + \kappa_2 \partial_y P) + \rho_{\star}^2 [\psi + \psi_0, \nabla_{\perp}^2 \psi] + \chi_{\perp} \nabla_{\perp}^2 P$ $\partial_t \psi = [\psi + \psi_0, \phi - P] - \partial_x P_0 \partial_y \psi + \eta \nabla_{\perp}^2 \psi + \eta C_b \partial_x P$

III. NL growth of NTM ?

Linear Spectrum



 Δ^{\prime} Δ^{\prime} Δ^{\prime} Δ^{\prime} Linear spectrum is stable with respect with tearing instability

=> No island

- Stable large scales modes
- Small scales turbulence driven by interchange instability
- Bootstrap current has a week effect on the linear spectrum
- => What 's about non-linear dynamics ?

III. NL amplification of TDMI by bootstrap current



Self-consistent generation of NTM from TDMI

1. TDMI formation => Seeding regime

2. NL growth of NTM => Amplification (by bootstrap current) regime [M. Muraglia et al, NF (2017)]

IV. Model in 3D cylindrical geometry

In 2D slab geometry, turbulence and magnetic island are located around the same resonant surface => Islands everywhere in tokamaks ?

 In fusion devices, no systematic overlap between turbulence area and magnetic island.



1 M.F.F Nave, et al. Nuc. Fus. 43 (2003) 2 A. Isayama, et al. Plasma and Fusion Research 8 (2013)

<u>3D cylindrical Reduced MHD model</u> :

$$\begin{aligned} \partial_t \tilde{\psi} &= \nabla_{\parallel} \tilde{\phi} - \nabla_{\parallel} \left(P_{eq} + \tilde{p} \right) + \eta \tilde{j} \\ \partial_t \tilde{\omega} &+ \left\{ \tilde{\phi}, \tilde{\omega} \right\} = \nabla_{\parallel} \left(J_{eq} + \tilde{j} \right) - \frac{\kappa_1}{r} \partial_{\theta} \tilde{p} + v \triangle_{\perp} \tilde{\omega} \\ \partial_t \tilde{p} &= \left\{ P_{eq} + \tilde{p}, \tilde{\phi} \right\} + \rho^{*2} \left\{ \Psi_{eq} + \tilde{\psi}, J_{eq} + \tilde{j} \right\} + \chi_{\perp} \triangle_{\perp} \tilde{p} \end{aligned}$$

A simple model including interchange and Tearing instabilities.

IV. Numerical set-up and Linear stage





 Edge turbulence level is controlled through pressure equilibrium gradient.

The q=2 surface is not located in an interchange area.

Parameters choosen such as q=2 surface is stable with respect to the tearing instability.

- No tearing mode develops at q=2.
- Interchange modes develop at the edge.

<u>Question</u> : Can we get a NL « spreading » of turbulence beating and generation of magnetic island at q=2 surface ?

IV. Nonlinear generation of q=2 island



The dominant mode, in nonlinear phase, is located at the lowest rational surface in the whole box : q=2. It is in a stable zone (with restect to interchange and tearing instability) and produces a (2, 1) magnetic island.

<u>Question</u> : What is the island generation mechanisms ?

IV. Beating of dominant unstable modes

Nonlinear beating rules :

1. The modes beat if they overlap.

The beating is effcient if the resulting mode is resonant at its birth location.



Then, how to explain the growth of the mode (2, 1)?

IV. Coeherent and non local beating

The beating mechanisms produces modes with large radial structure (5, 2), (7, 3) and (9, 4) in the QL phase and remains efficient in the whole NL phase.

The beating of such modes generates (2, 1) but only at the tail of the eigen function, at q=2.

(9,4) + (7,3) = (2,1)(7,3) + (5,2) = (2,1)





[A. Poyé et al, POP 22 (2015)]

IV. Toroidal source-driven model



Simple flux-driven model for toroidal interchange $\partial_t \omega + [\phi, \omega] = -\kappa_1 \mathcal{G}P + \nabla_{\parallel} J + \nu \Delta_{\perp} \omega$ $\partial_t P + [\phi, P] = \kappa_2 \mathcal{G}\phi + \kappa_2^P \mathcal{G}P + \rho_{\star}^2 [\psi, J] + \chi \Delta_{\perp} P + S_P$ $\partial_t \psi = \nabla_{\parallel} (\phi - P) + \eta j_{\parallel}$

where

$$\mathcal{G} = \sin\theta\partial_r + \frac{1}{r}\cos\theta\partial_\theta$$

is the curvature operator.



Source term :

- constant and controlled profile throughout saturation phase

- allows consistent dynamics close to marginal stabiltiy

IV. Multiple coupling possibilities

 $\tilde{\phi}(r, \theta_{\overline{n}}/\varphi = 0)$



0

3π/2





$$\mathcal{G}f = \sin\theta\partial_r f + \frac{1}{r}\cos\theta\partial_\theta f$$

Additional $m \leftrightarrow m\pm 1$ coupling

Mode couplings :

- Local nonlinear (conserves parity)
- Nonlocal nonlinear
- Linear toroidal



IV. Toroidal coupling



(2,1) damping (3,1) damping –

1.40

1.45

1.50 1e4

1.35

10⁻¹³ 1.20

1.25

1.30

Direct coupling ? (3,1) in the turbulent zone (2,1) in the stable zone ?

> Kill artificially (2,1) mode Let it recover

- Suppressing (3,1) mode
- Leaving (3,1) mode

Same recovery in both cases

Role of (3,1) mode negligible

IV. Toroidal coupling



IV. Spectrum dynamics



Increasing power source (feeding more energy to turbulence)

Remote island spectrum evolves towards small mode numbers

Direct toroidal coupling of medium-scale modes

Nonlinear coupling of medium scale modes

all modes



IV. Spectrum dynamics



Special case : tearing-like modes slightly unstable

 \rightarrow Dominated by medium-scale tearing (5,2), (7,3)

Final state goes from n=3 to n=1 with increasing source

Intermediate final values \rightarrow oscillating/chaotic regimes

Integer final values \rightarrow steady-state regime (with lower maximum island size)

Also : generation mechanism *independent on underlying instability*



 $\Psi_0 + \tilde{\Psi}(r, \theta, 0) - \frac{r^2 2\pi}{2q_s L_z}$ 2π 3π/2 θ π π/2 0 0.54 0.58 0.6 0.62 0.66 0.56 0.64 0.68 0.7 Circular concentric flux surfaces are not a toroidal equilibrium state

$$\begin{aligned} \partial_t \omega + [\phi, \omega] &= -\kappa_1 \mathcal{G}P + \nabla_{\parallel} J + \nu \Delta_{\perp} \omega \\ \partial_t P + [\phi, P] &= \kappa_2 \mathcal{G}\phi + \kappa_2^P \mathcal{G}P + \rho_{\star}^2 [\psi, J] + \chi \Delta_{\perp} P + S_P \\ \partial_t \psi &= \nabla_{\parallel} (\phi - P) + \eta j_{\parallel} \end{aligned}$$

→ Generation of an axisymmetric perturbation of ψ and Φ from pressure gradient

Mostly irrelevant (~static, low gradients)

However, must be taken into account when computing helical flux $r^2 2\pi$

$$\chi = \psi_0 + \psi - \frac{7}{2q_s} \frac{2\pi}{L_z}$$

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 $\Gamma_{\rm S}$

Removing axisymmetric (equilibrium) perturbation allows to recover

- overall island structure
- X-points, O-points phases
- Island width (order of magnitude)

Or : only keep resonant perturbation

However that's still a 2D diagnostic

- Each φ-plane gives a different island
- KAM islands in stochastic region

Island size from Poincaré section ?



Removing axisymmetric (equilibrium) perturbation allows to recover

- overall island structure
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Island size from Poincaré section?

Island size from Poincaré map :

3 kind of field lines :

- « regular » field lines
- field lines « trapped » inside island
- stochastic region

Field lines inside an island don't have θ values close to that of X point(s) \rightarrow gap in θ values

Island size is approached by the radial width of the largest field line trapped within an island

Average r value \rightarrow resonant surface

Only free parameter : gap threshold





Island size proportional to power source

Weaker dependence after a threshold corresponding to turbulent stochastic zone spreading

- \rightarrow More efficient transport for a given
- power source & fluctuation level ? No
- \rightarrow Stochasticity limits island size ?

Not in 2D case !

Limited dependance on distance

Same order of magnitude for local and remote coupling

Robust mechanism



Field lines in the stochastic region near an island can have a quasi-regular structure

- over >2500 returns to Poincaré plane
- beware numerical effects !

Probably not topological field line may ultimately escape

However, relevant for transport / NTM

Correctly catalogued as island field line.

 $\rightarrow\,$ island size diagnostic includes green region



V. Conclusions

A basic RMHD model has been used to investigate the interaction between small scale interchange turbulence and magnetic islands when tearing instability is marginally stable.

In 2D slab geometry, it has been shown that small scale turbulence leads to the NL generation of seed magnetic islands thanks to a interchange modes beating. The island and the turbulence are located around the same resonant surface.

Bootstrap current has been added in the 2D slab model. Self-consitent generation of NTMs from turbulence driven seed island.

In 3D geometry, nonlinear simulations have been performed in cases where interchange instability plays at the edge whereas the q=2 surface is located in an inner area. At the edge, efficient nonlinear beating of interchange modes generates large radial tearing modes which overlap the q=2 surface. Thus a (2, 1) magnetic island is generated.

Numerical tool : AMON code

Semi-spectral code :

Radial direction : finite difference

Poloidal and axial (for 3D) directions : spectral

<u>Resolution used for this study</u> : 1024 points in the radial direction and 256 poloidal modes

Temporal scheme :

Runge-Kutta 4

Boundary conditions :

Radial direction : 0 at the boundaries Poloidal and axial directions : periodic

Nonlinear terms :

Quadratic terms conservation

Numerical tool : Performances code

1024*512, 2D simulations
 with 2 fields, 500000 iterations
 on Nestor (8 procs/node) :

256*256*128, 3D simulations
 with 3 fields, 500000 iterations
 on Juelich (8 procs/node) :



Beating of interchange modes

•
$$\Delta' = [rac{{
m d} \ln \psi_{
m ideal}}{{
m d} {
m x}}]_{{
m 0}^-}^{{
m 0}^+} \in [-0.5, 4.5]$$

Tearing marginally stable/unstable

> Interchange instability growth rate depends on the equilibrium and therefore is linked to Δ '

> Interchange is enhanced when Δ' becomes negative $(d\gamma_{\star}/d\Delta' < 0)$





 Beating of small scales interchange modes drive large scale modes:

 $2\gamma_{\star} > \gamma_1^L$: Interchange driven island $2\gamma_{\star} < \gamma_1^L$: Tearing driven island

III. Perturbed current profile



Perturbed current profile is not enough to destabilize islands.

Island generation is not due to tearing + anomalous resistivity

III. Profile stiffness

