



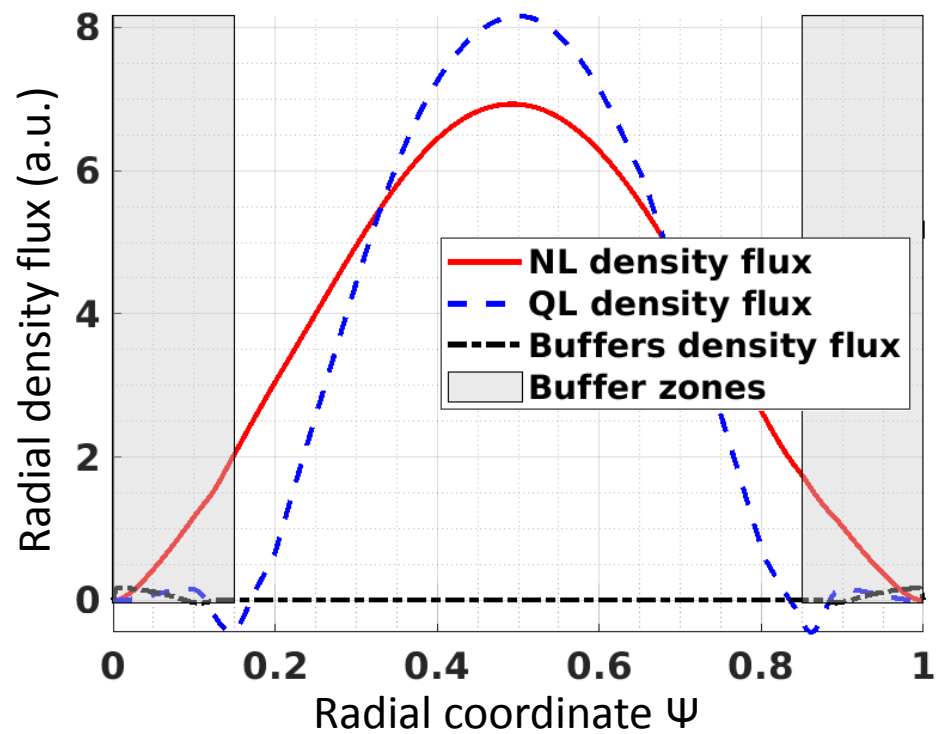
Anatomy of turbulent transport in energy-space

J. Médina, M. Lesur, E. Gravier

Collaborators: X. Garbet, T. Cartier-Michaud, T. Drouot,
M. Idouakass, Y. Kosuga, T. Réveillé,
P. Bertrand, P.H. Diamond

Fundings: French Research Federation for Magnetic Confined Fusion

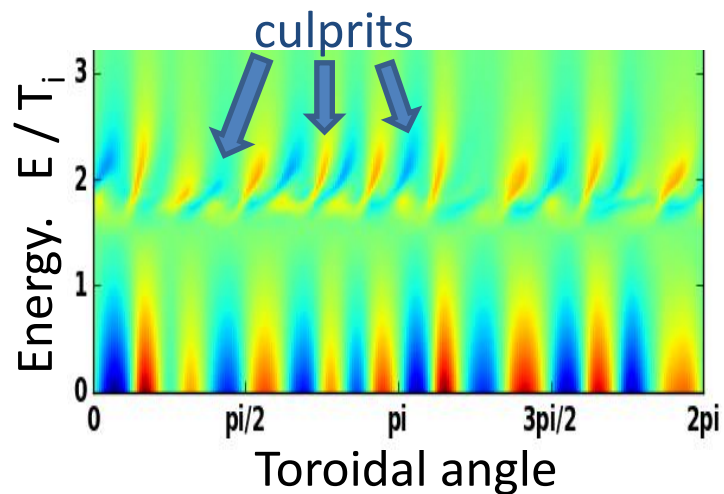
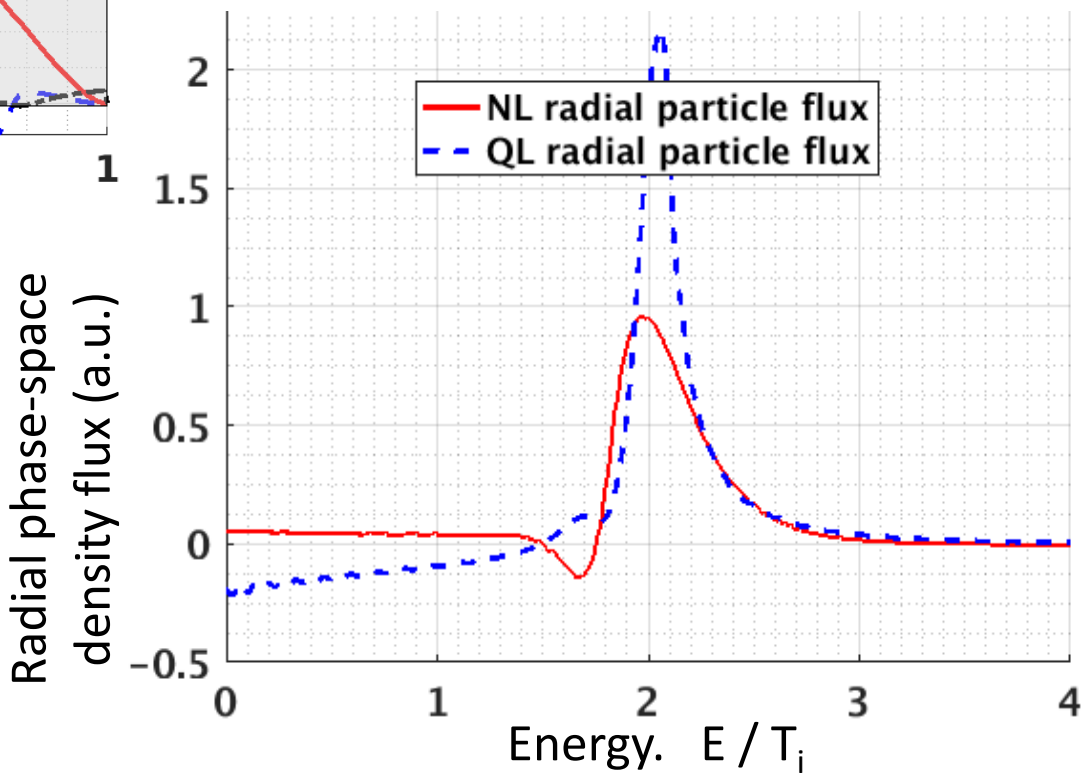
Conclusion: quasilinear theory works!...?...?!...



GK simulation of turbulent transport in tokamak core

← Good prediction

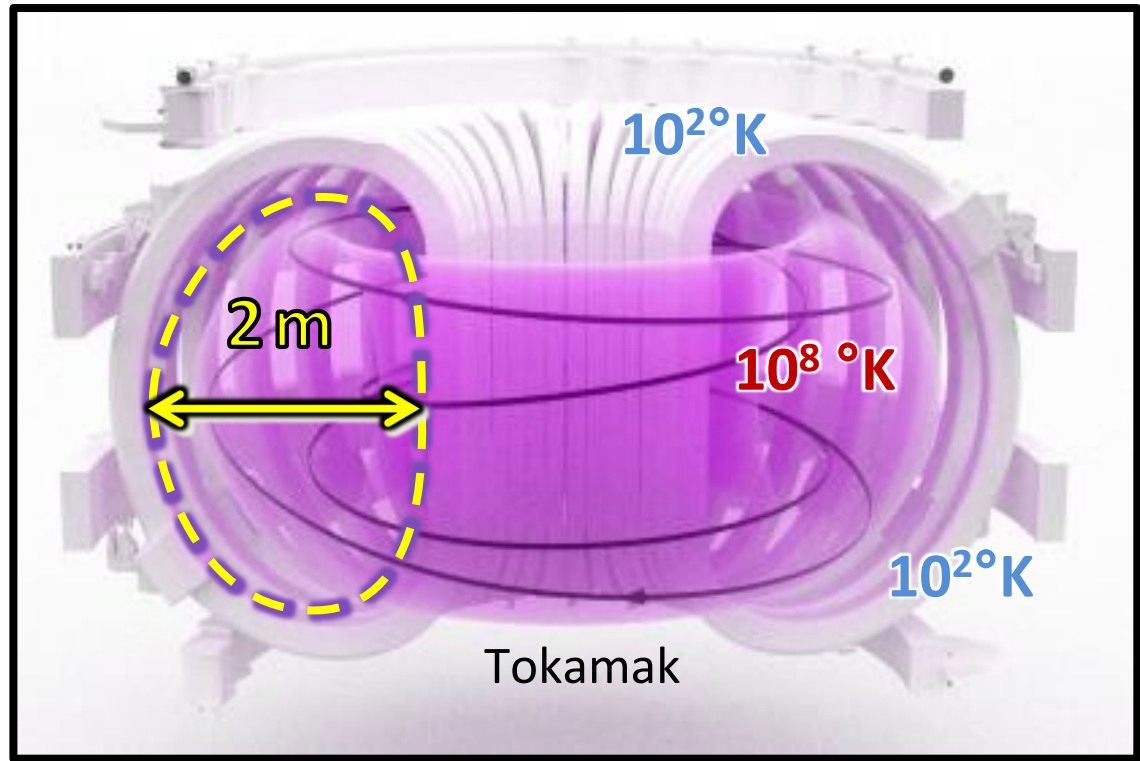
for the wrong reasons



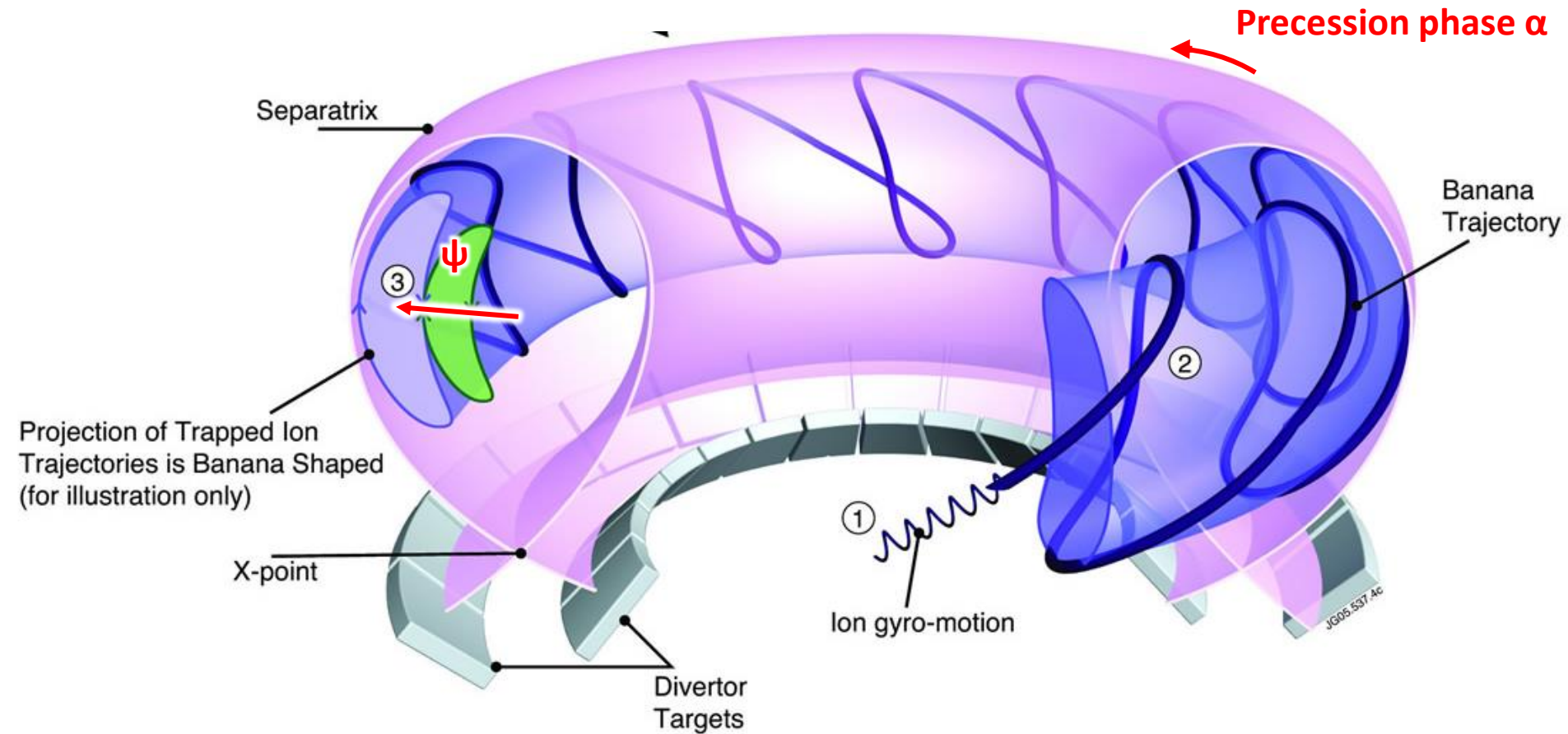
Context: Turbulence in the core of fusion plasmas

Core turbulence

- Microscopic (mm) waves
- Driven by pressure gradients
- Dominated by Ion-Temperature-Gradient (ITG) mode and Trapped-Electron-Mode (TEM)
- Collisional mean free path (km) $\gg \lambda$
- In this work: electrostatic, collisionless
- Transport particles and energy between core and edge via **ExB** drift
 - \Rightarrow Degrade confinement
 - \Rightarrow May be harnessed



Magnetically trapped particles (bananas)



Deeply trapped particles described by 3 variables: α , ψ , E

Radial coordinate ψ Kinetic energy E

Approach: bounce-averaged gyrokinetics

Kinetics:

6D Phase-space (3D + 3V)



$$\omega \ll \omega_{\text{gyro}}$$

Gyrokinetics:

4D Phase-space + 1 parameter



$$\omega \ll \omega_{\text{bounce}} \ll \omega_{\text{gyro}}$$

Gyrobounce
gyrokinetics:

2D Phase-space + 2 parameters



$$\text{Pitch-angle} = \pi/2$$

2D Phase-space + 1 parameter

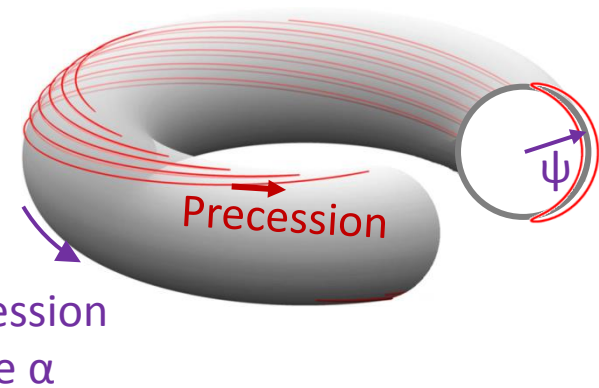
(α , ψ)

(Kinetic energy E)

Main limitation: $\omega \sim \omega_{\text{precession}}$

⇒ Trapped-Ion Mode (TIM) and
Trapped-Electron Mode (TEM)

⇒ Between DNS and reduced model



Model equations

Vlasov equation: $\frac{\partial f}{\partial t} + \underbrace{[J_0 \phi, f]_{\alpha, \psi}}_{\mathbf{E} \times \mathbf{B} \text{ motion}} + \underbrace{E \Omega_d \frac{\partial f}{\partial \alpha}}_{\text{Precession}} = 0$ $\dot{\alpha} = E \Omega_d + \frac{\partial \bar{\phi}}{\partial \psi}$
 $\dot{\psi} = -\frac{\partial \bar{\phi}}{\partial \alpha}$

Quasi-neutrality:

$$\underbrace{C_1 [\phi - \langle \phi \rangle + \mathcal{F}^{-1}(i \delta_m \hat{\phi}_m)]}_{\delta n_{\text{passing}}} - \underbrace{C_2 \bar{\Delta} \phi}_{\text{Polarisation}} = \underbrace{\frac{2}{\sqrt{\pi}} \int_0^\infty J_0(E) f \sqrt{E} dE}_{\bar{\bar{n}}_{\text{trapped}}} - 1$$

$\sim (\delta_b k_r)^2$ GK average ≈ 1
 \downarrow \downarrow

Depret '00

Tagger, Laval, Pellat '77

Sarazin '05

Biglari, Diamond, Terry '88

Darmet '08

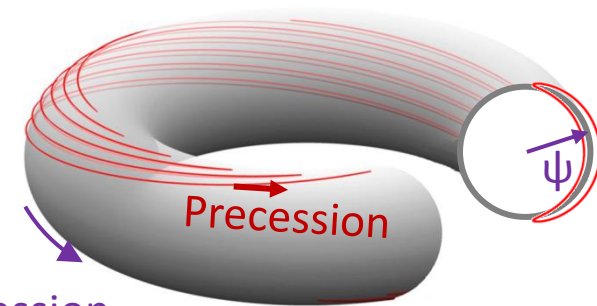
\Rightarrow TERESA code

Cartier-Michaud '13

Drouot '14

- Global, full-f, fixed-gradient (flux-driven soon)
- Collisionless, electrostatic
- Quasi-adiabatic passing particles

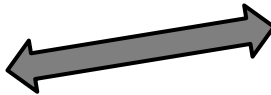
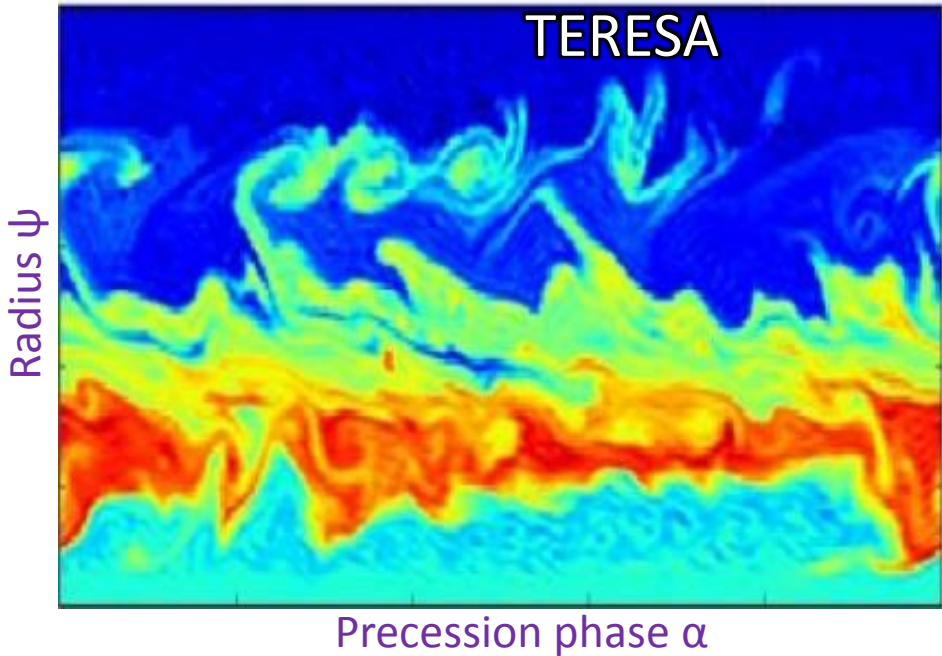
Lesur '17



Precession
phase α

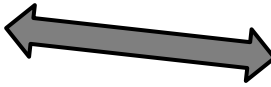
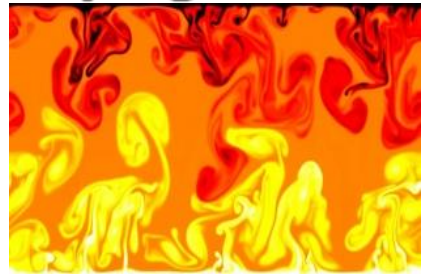
TERESA explores kinetic effects in tokamak turbulence

Kinetic version of classic instabilities



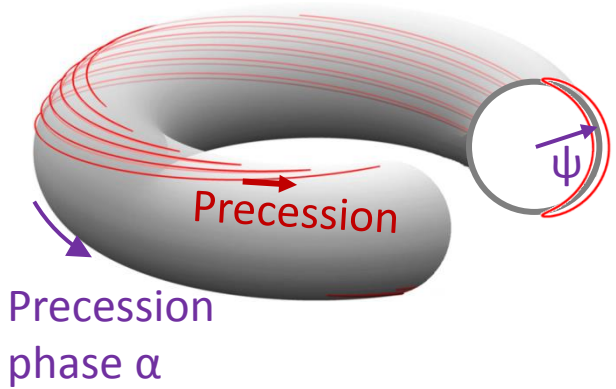
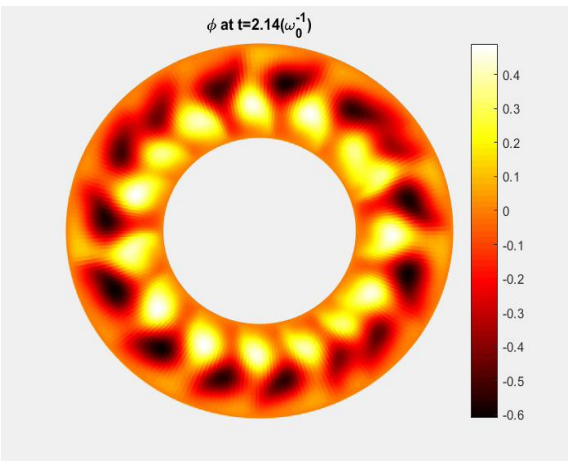
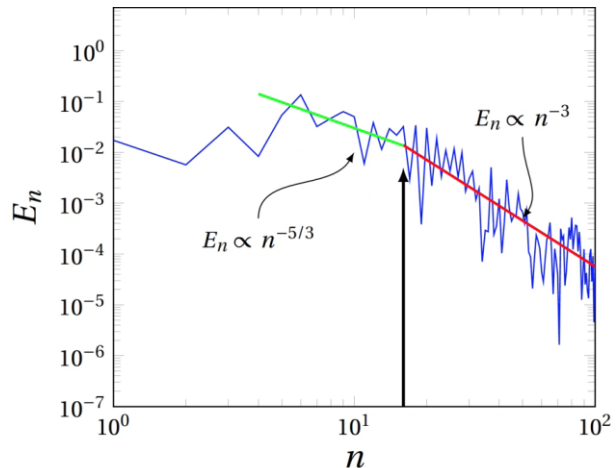
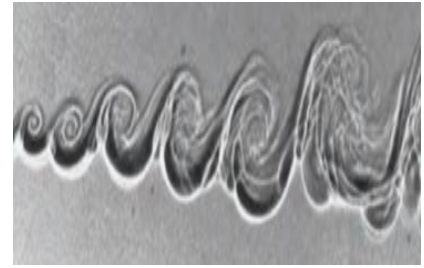
Cartier-Michaud '14

Rayleigh-Bénard



Palermo '15

Kelvin-Helmholtz



Derivation of quasilinear theory

Vlasov in Fourier space

$$\frac{\partial \hat{f}_m}{\partial t} + \underbrace{im \frac{E\Omega_d}{Z} \hat{f}_m}_{\text{Free streaming (precession)}} - \underbrace{im \hat{\phi}_m \frac{\partial \hat{f}_0}{\partial \psi}}_{\text{Linear drive}} = - \sum_l il \hat{f}_l \frac{\partial \hat{\phi}_{m-l}}{\partial \psi} + \underbrace{\sum_{l \neq 0} i(m-l) \hat{\phi}_{m-l} \frac{\partial \hat{f}_l}{\partial \psi}}_{\text{Nonlinear terms}}$$

Radial flux (phase dynamics)

$$m = 0. \quad \frac{\partial \langle f \rangle}{\partial t} + \sum_l il \frac{\partial (\hat{f}_l \hat{\phi}_l^*)}{\partial \psi} = 0. \quad \Rightarrow \text{Evolution of mean fields}$$

$$m \neq 0 \quad \frac{\partial \hat{f}_m}{\partial t} + im \frac{E\Omega_d}{Z} \hat{f}_m - im \hat{\phi}_m \frac{\partial \langle f \rangle}{\partial \psi} + im \hat{f}_m \frac{\partial \hat{\phi}_0}{\partial \psi} = 0$$

Neglecting all nonlinearities except for $l = m$

Derivation of quasilinear theory (2)

$$m \neq 0 \quad \frac{\partial \hat{f}_m}{\partial t} + im \frac{E \Omega_d}{Z} \hat{f}_m - im \hat{\phi}_m \frac{\partial \langle f \rangle}{\partial \psi} + im \hat{f}_m \frac{\partial \hat{\phi}_0}{\partial \psi} = 0$$

$L \hat{f}_m$

⇒ Linear response

$$\hat{f}_m(\psi, E, t) = \int_0^t e^{i\omega_{R,m}(s-t)} im \hat{\phi}_m(\psi, E, s) \frac{\partial \langle f \rangle}{\partial \psi}(\psi, E, s) ds.$$

$$\omega_{R,m}(\psi, E, t) = m \left(\frac{\Omega_d E}{Z} + \frac{\partial \hat{\phi}_0}{\partial \psi} \right)$$

Doppler-shift
by zonal flow

Vedenov, Velikov, Sagdeev '61
Drummond & Pines '62
Sagdeev & Galeev '69

Derivation of quasilinear theory (3)

Linear theory

$$\hat{\phi}_m(\psi, E, t) = \hat{\phi}_m(\psi, E, 0) \exp \left[\int_0^t (-i\omega_m + \gamma_m(t')) dt' \right]$$

$$\frac{\partial \langle f \rangle}{\partial t} = \frac{\partial}{\partial \psi} \left[\sum_l l^2 |\hat{\phi}_l(\psi, E, t)|^2 \int_0^t e^{i(\omega_{R,l} - \omega_l)(s-t)} \frac{\partial \langle f \rangle}{\partial \psi} \Big|_s \right. \\ \left. \times \exp \left(\int_t^s \gamma(t') dt' \right) ds \right]$$

Phase-mixing for $t-s >$ growth, relaxation times

$$\Rightarrow \frac{\partial \langle f \rangle}{\partial t} = \frac{\partial}{\partial \psi} \left[\underbrace{D_{\text{QL}}}_{\Lambda_{\psi}^{\text{QL}}(\psi, E, t)} \frac{\partial \langle f \rangle}{\partial \psi} \right]$$

Vedenov, Velikov, Sagdeev '61

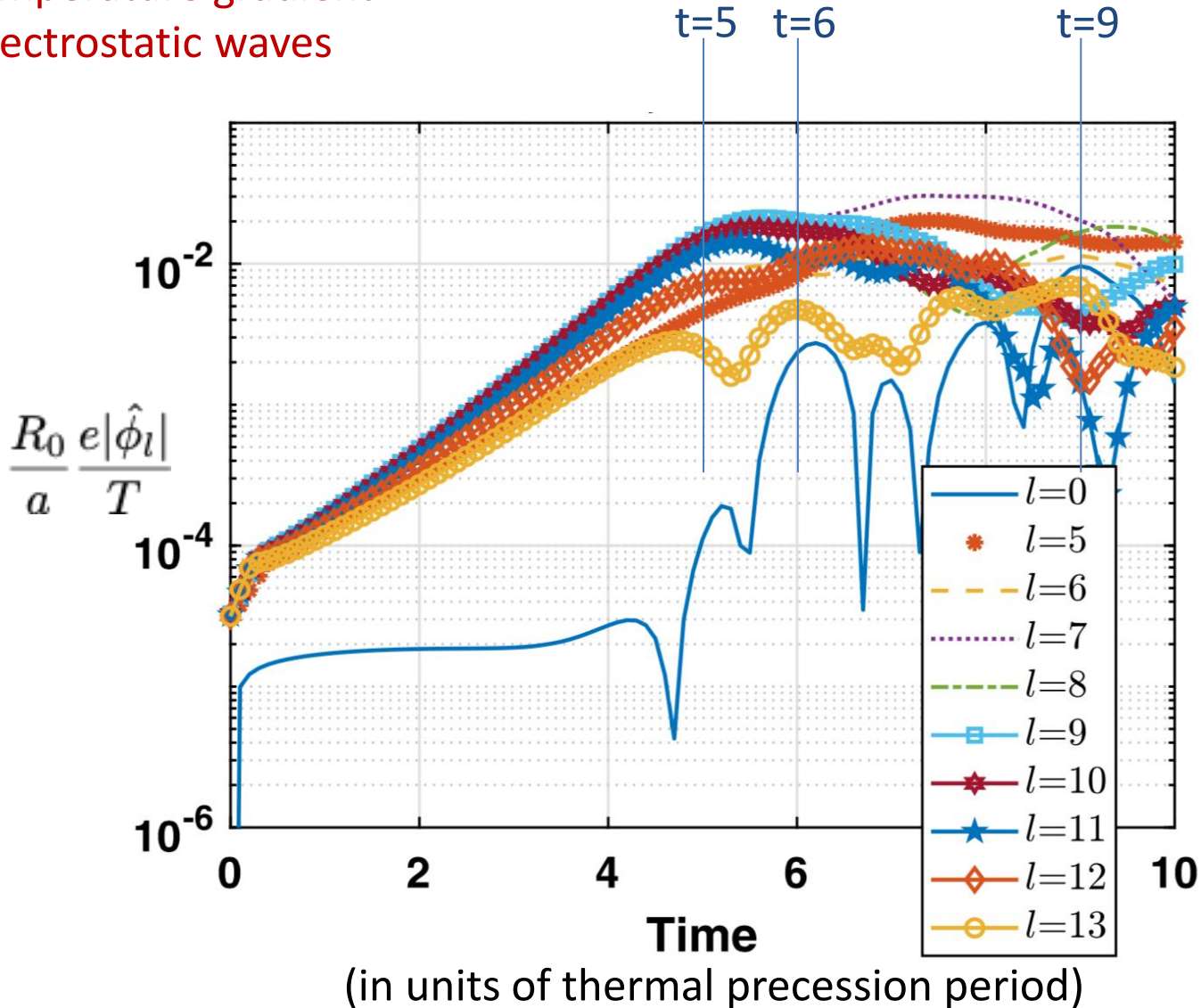
Drummond & Pines '62

Sagdeev & Galeev '69

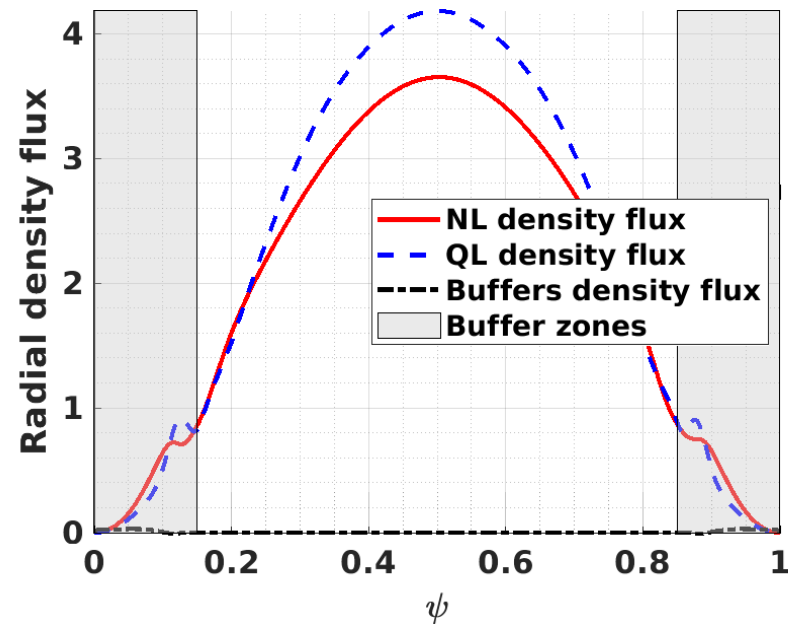
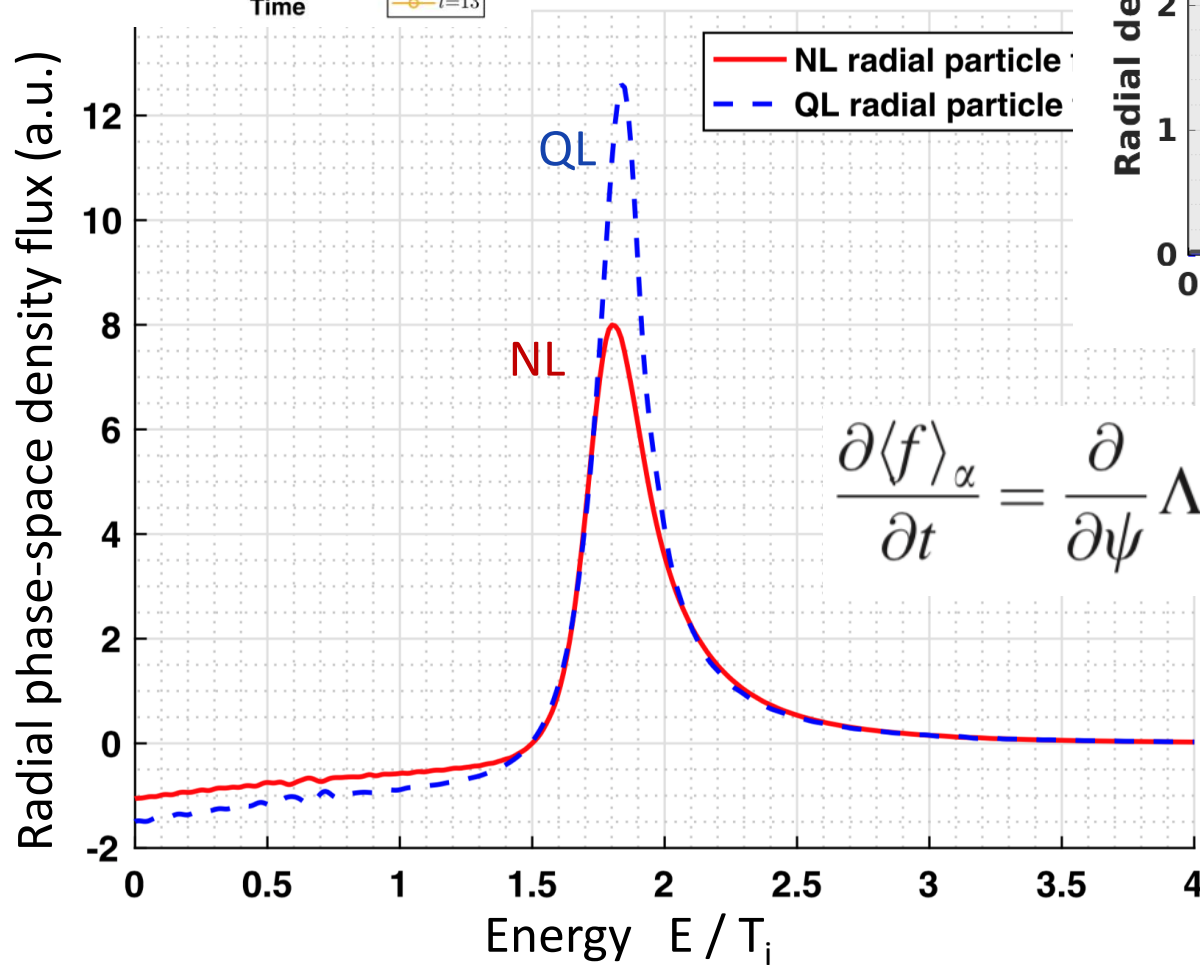
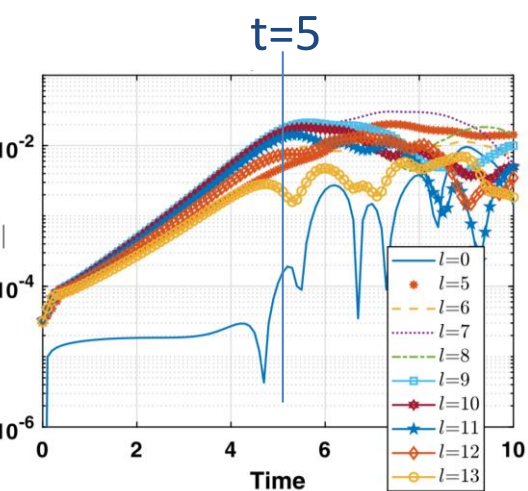
$$D_{\text{QL}}(\psi, E, t) = \sum_l l^2 |\hat{\phi}_l(\psi, E, t)|^2 \frac{1 - e^{-i(\omega_{R,l} - \omega_l)t - \gamma_l t}}{i(\omega_{R,l} - \omega_l) + \gamma_l}$$

Trapped-ion-driven turbulence

Initial temperature gradient
drives electrostatic waves



Anatomy of transport



$$\frac{\partial \langle f \rangle_\alpha}{\partial t} = \frac{\partial}{\partial \psi} \Lambda_\psi^{\text{NL}}(\psi, E, t)$$

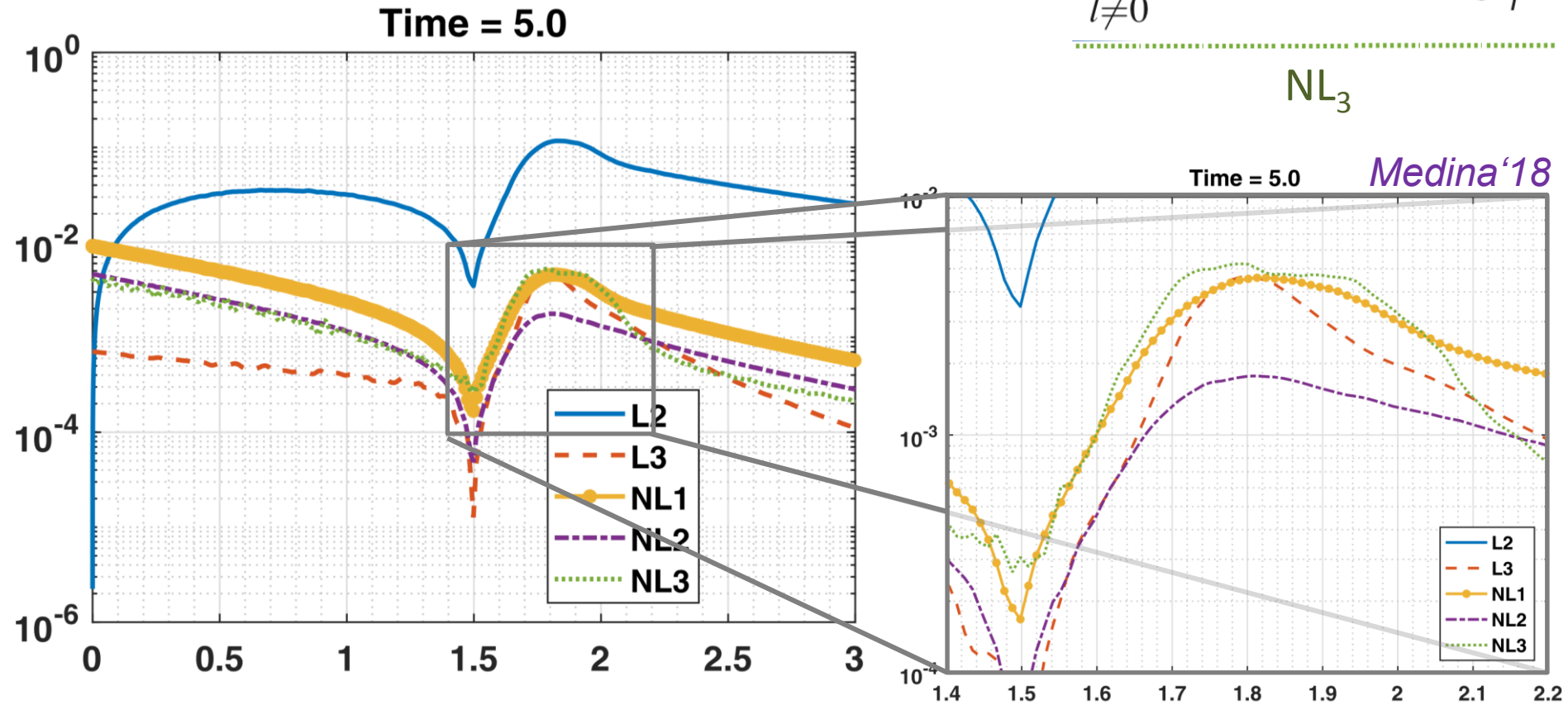
Radial phase-space density flux

Medina'18

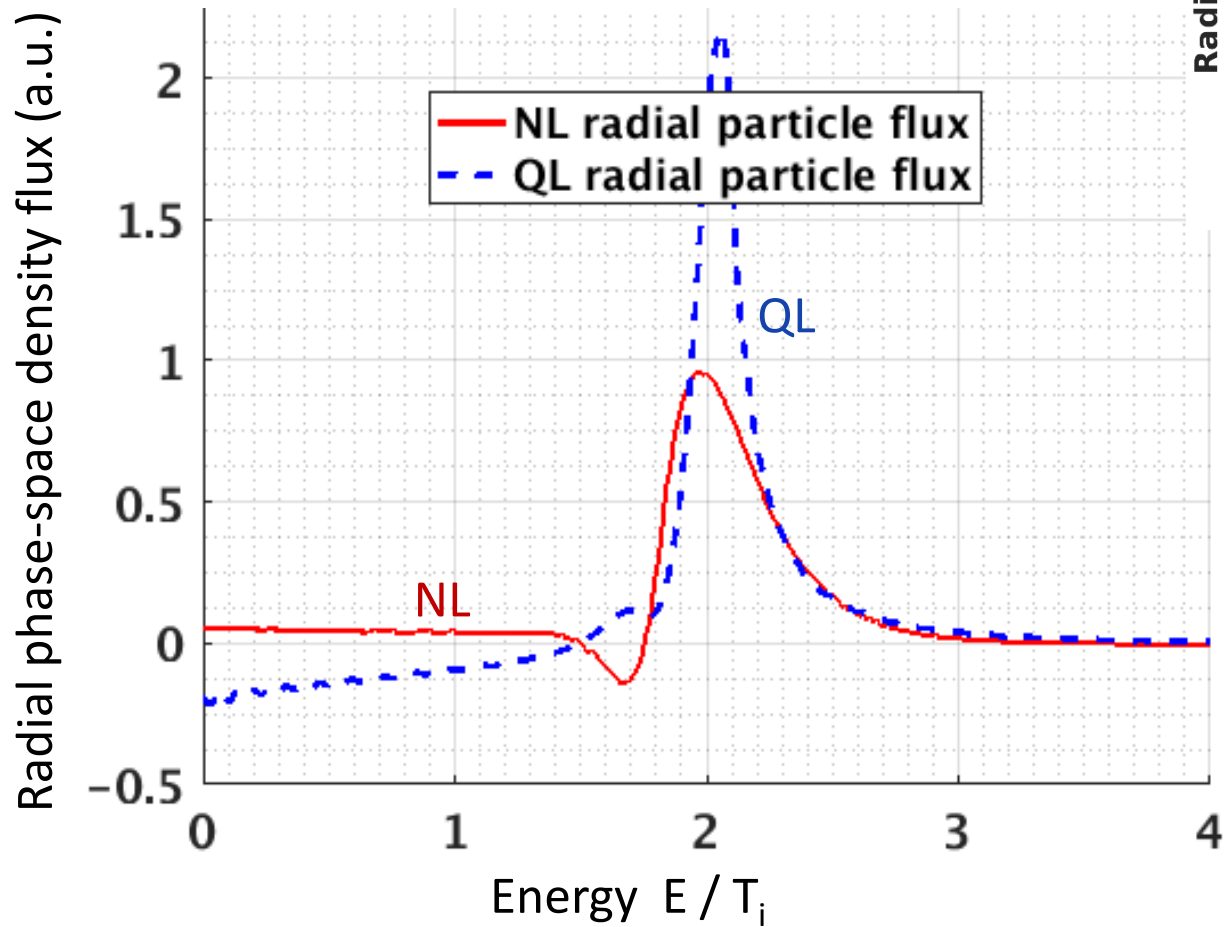
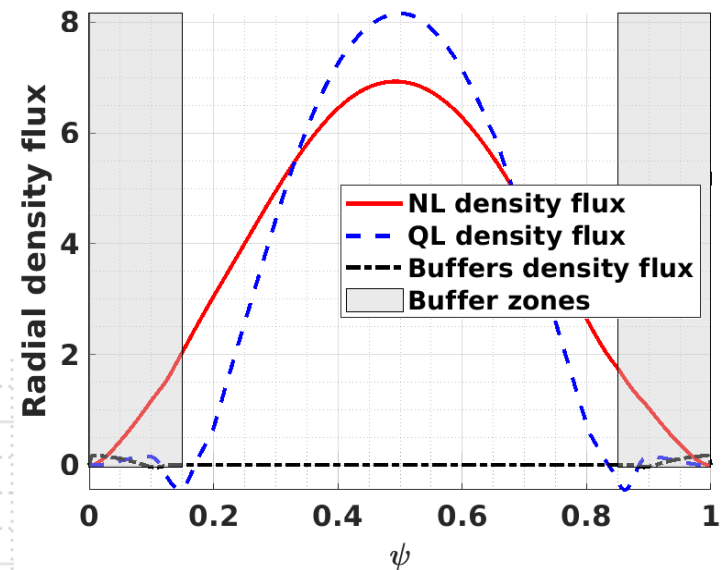
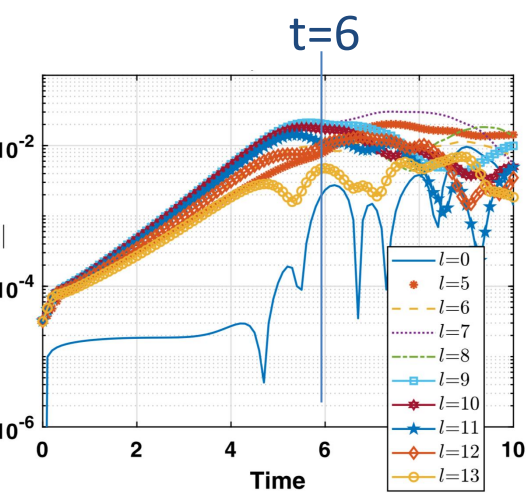
Neglected NL terms are important in magnitude

$$\underbrace{\frac{\partial \hat{f}_m}{\partial t}}_{L_1} + \underbrace{im \frac{E\Omega_d}{Z} \hat{f}_m}_{L_2} - \underbrace{im \hat{\phi}_m \frac{\partial \hat{f}_0}{\partial \psi}}_{L_3} = - \sum_l ilf_l \frac{\partial \hat{\phi}_{m-l}}{\partial \psi} + \underbrace{\sum_{l \neq 0} i(m-l) \hat{\phi}_{m-l} \frac{\partial \hat{f}_l}{\partial \psi}}_{NL_3}$$

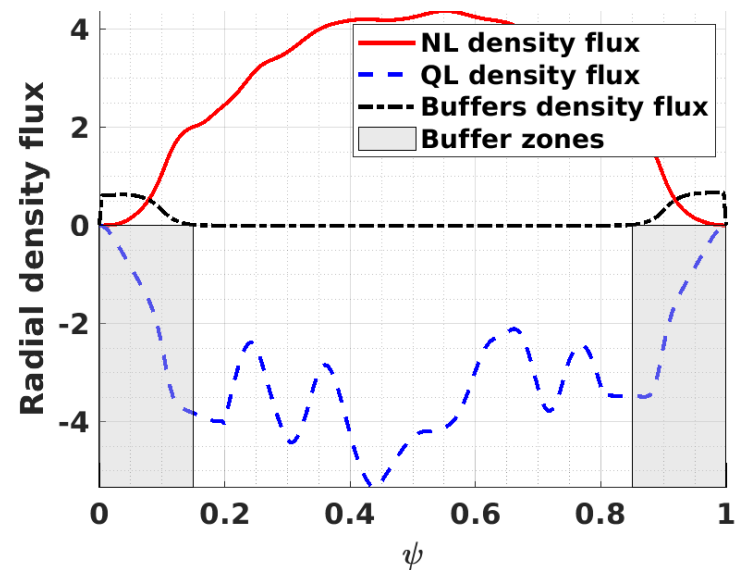
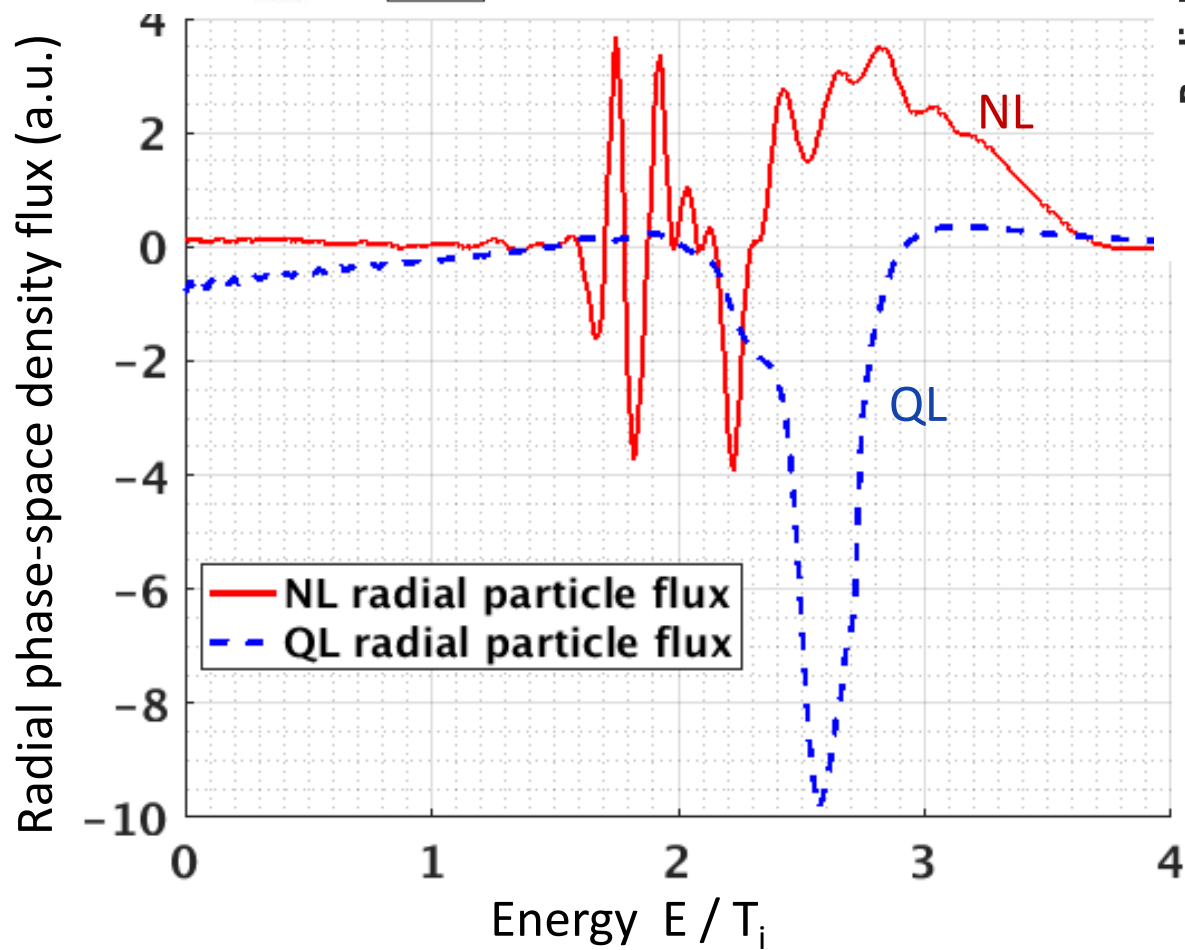
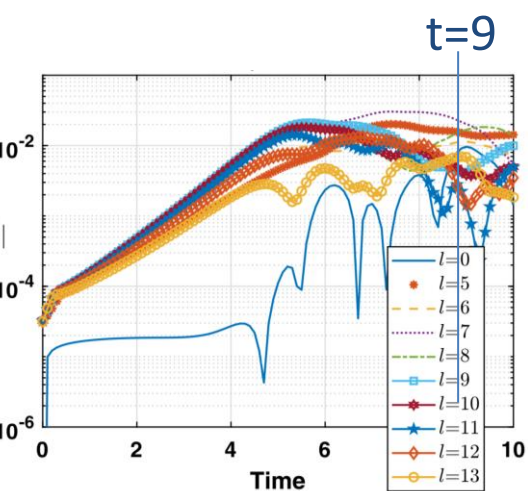
↗ **NL₁ (l=m)**
↘ **NL₂ (l≠m)**



Anatomy of transport (2)

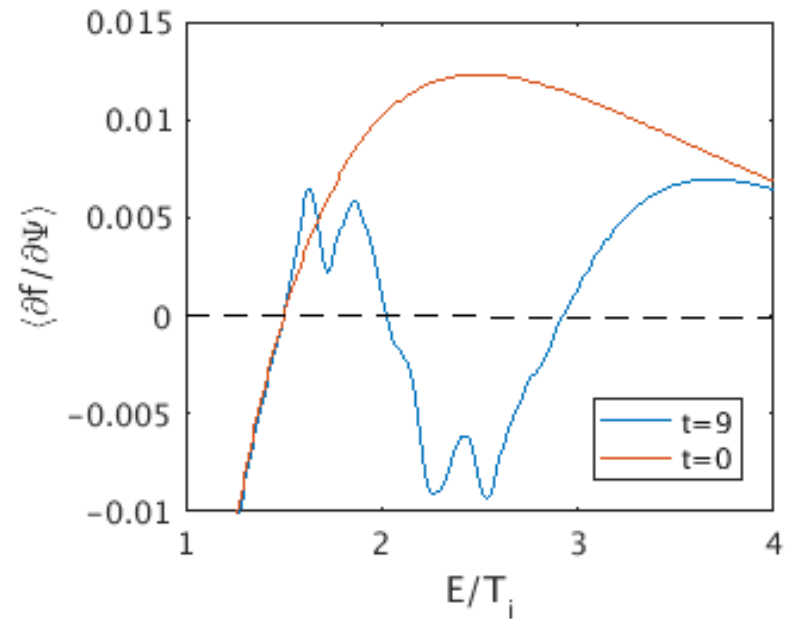
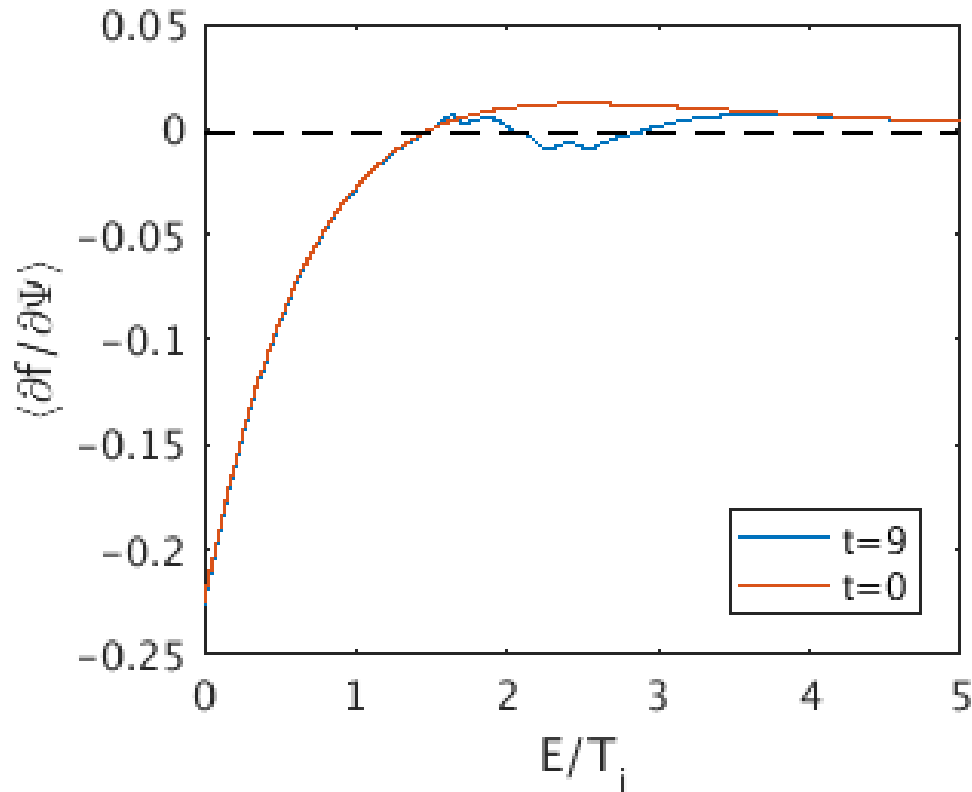


Anatomy of transport (3)

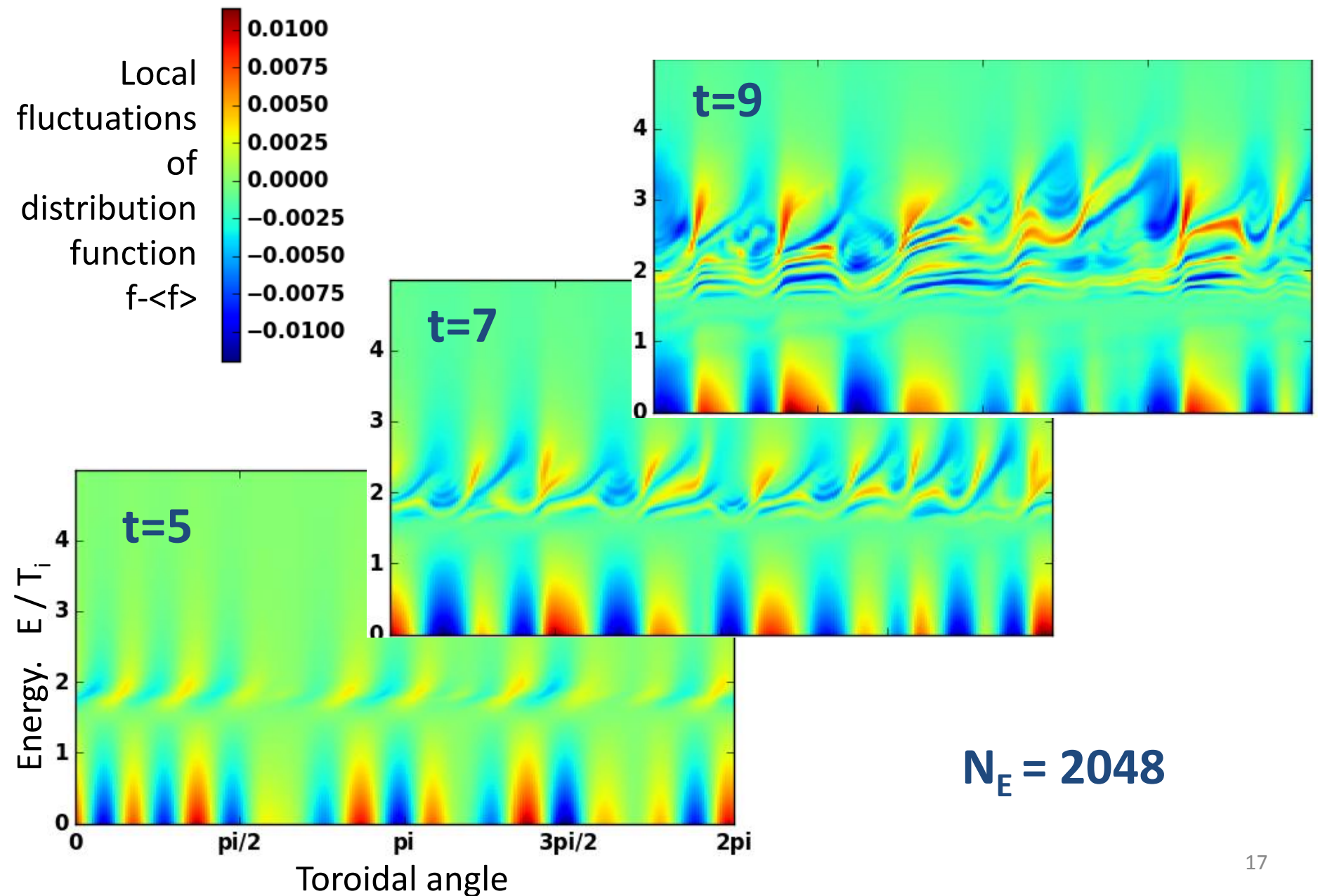


QL flux < 0 due to fine scales in energy-space

$$\Lambda_{\psi}^{\text{QL}}(\psi, E, t) = -D_{\text{QL}}(\psi, E, t) \frac{\partial \langle f \rangle_{\alpha}}{\partial \psi}(\psi, E, t)$$



Structures in phase-space



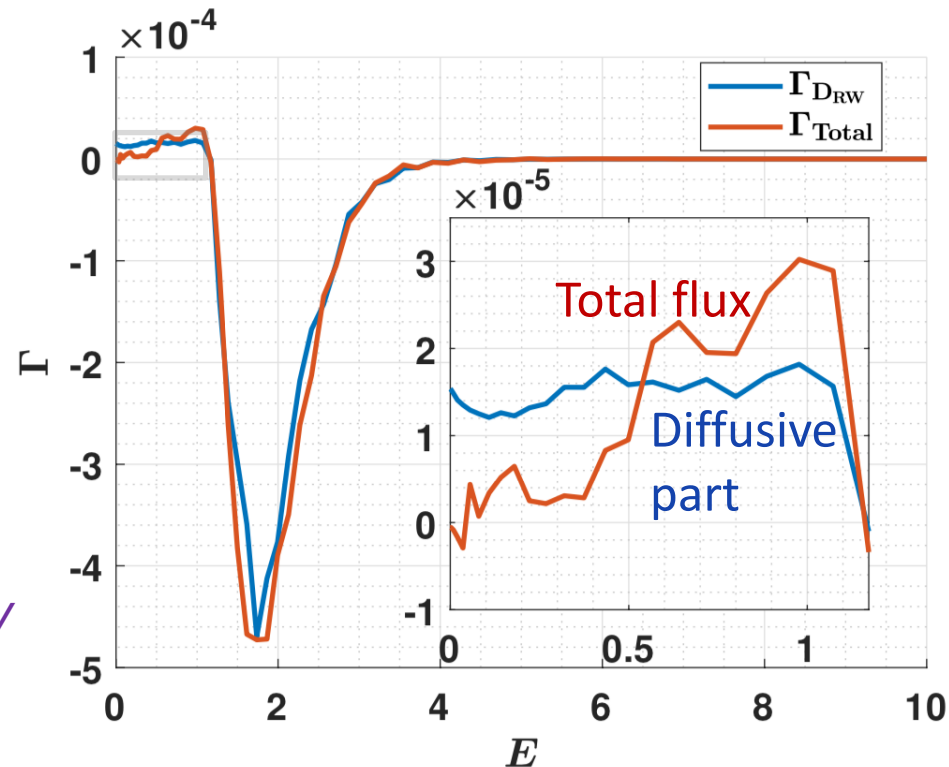
Summary

Collisionless turbulent transport is dominated by resonant processes

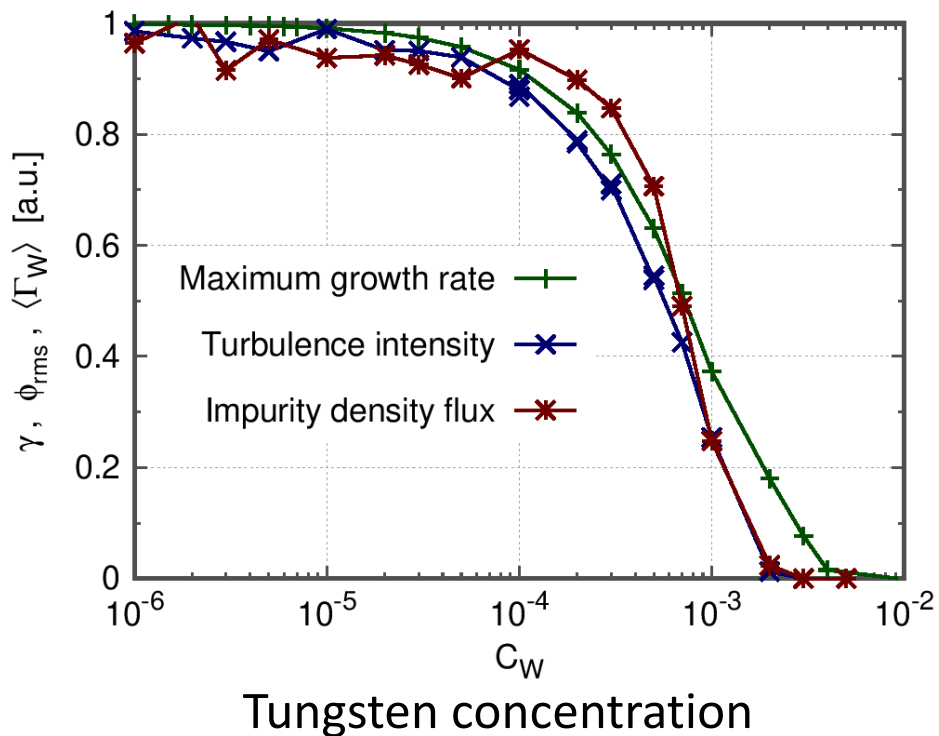
- Radial flux of particles dominated by a narrow peak around resonance
- Caveat: trapped-particle-driven only (strong resonances)
- Early nonlinear stage: diffusion in agreement with quasilinear theory
- Later stage: departure from quasilinear formalism due to small structures in energy-space

Perspectives

- Weak vs stronger turbulence?
- Diffusion vs convection vs trapping?
- Steady-state in a stirred system (flux-driven)?
- Collisions \longrightarrow *Artur Kryzhanovskyy*
- Beyond QL theory?

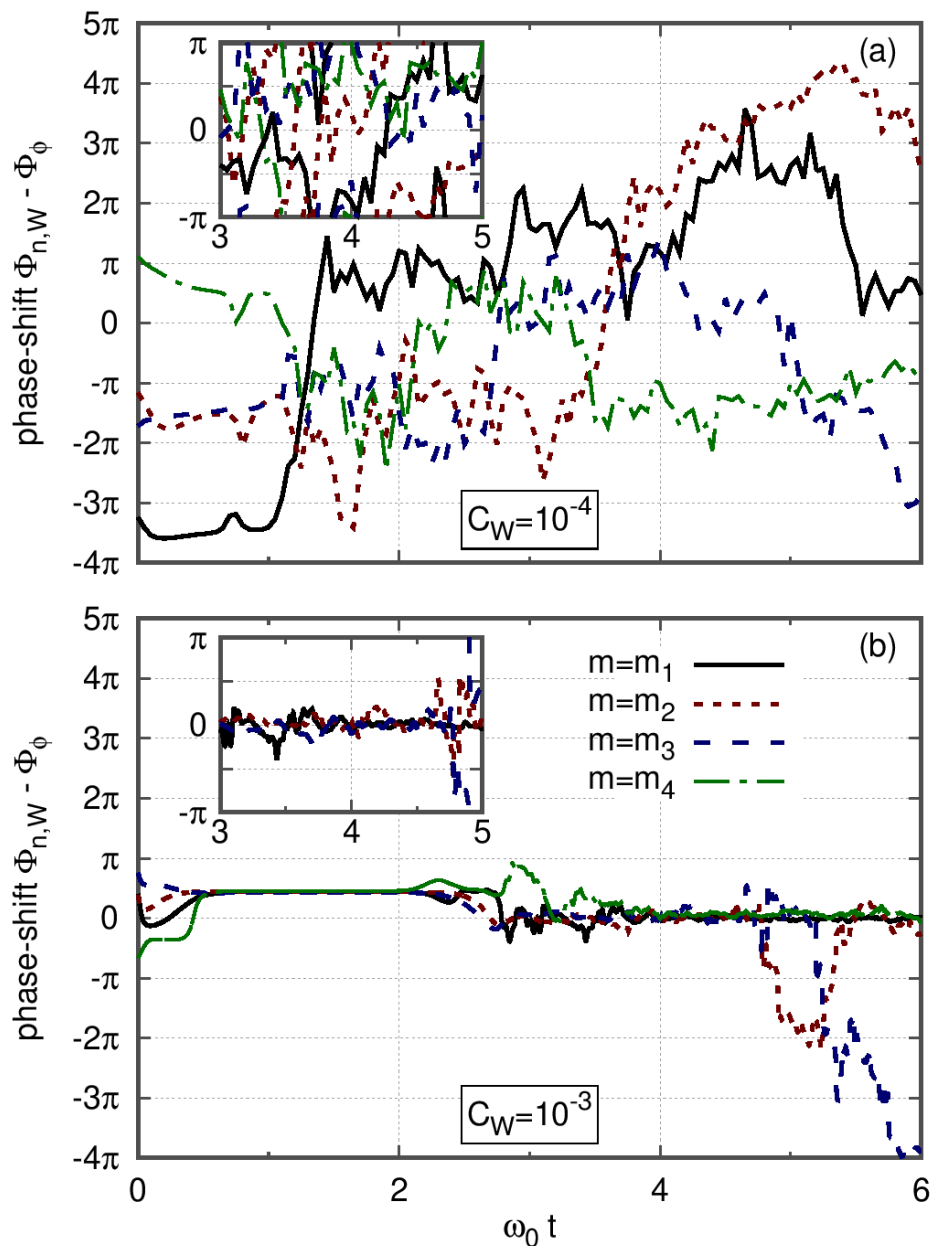


Bonus on phase dynamics: impurity transport

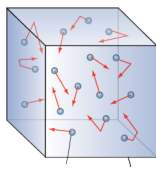


⇒ Quenching of impurity transport by synchronisation of potential fluctuations on impurity density fluctuations

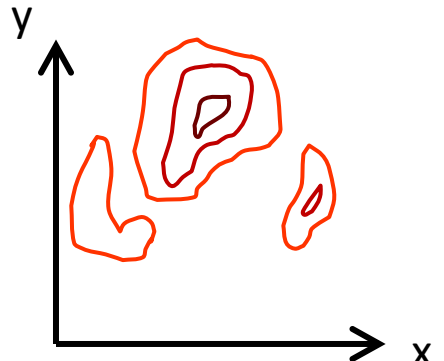
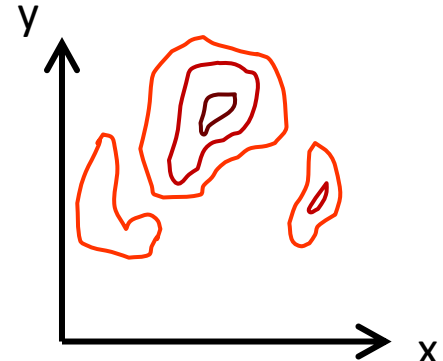
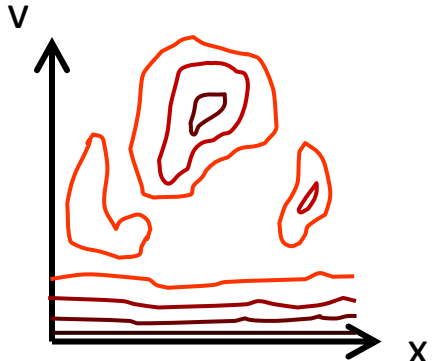
Threshold?



Additional material

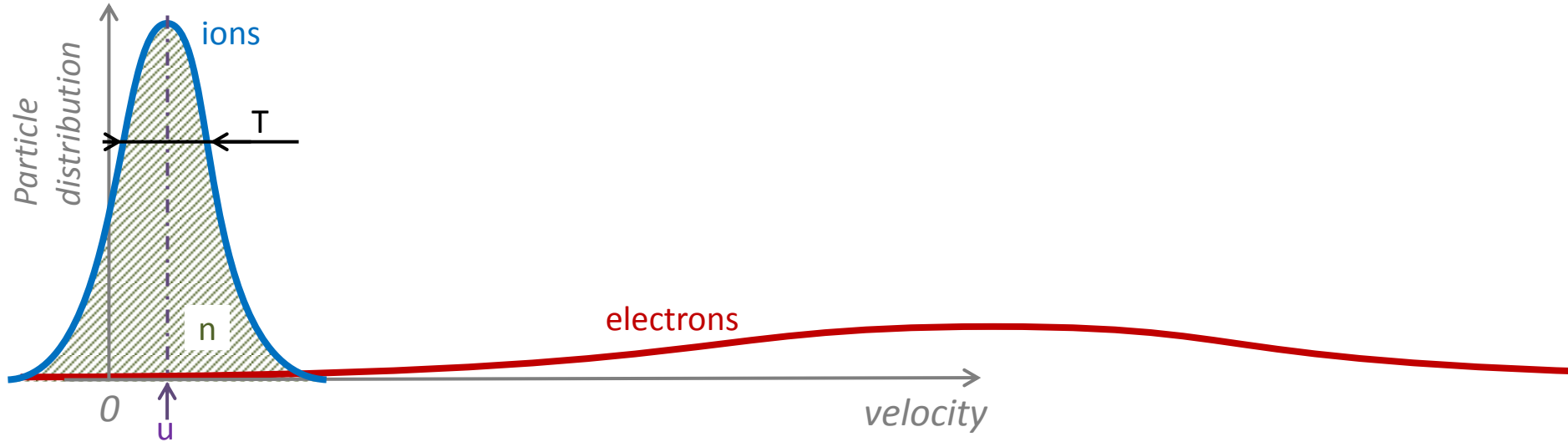


Kinetic theory

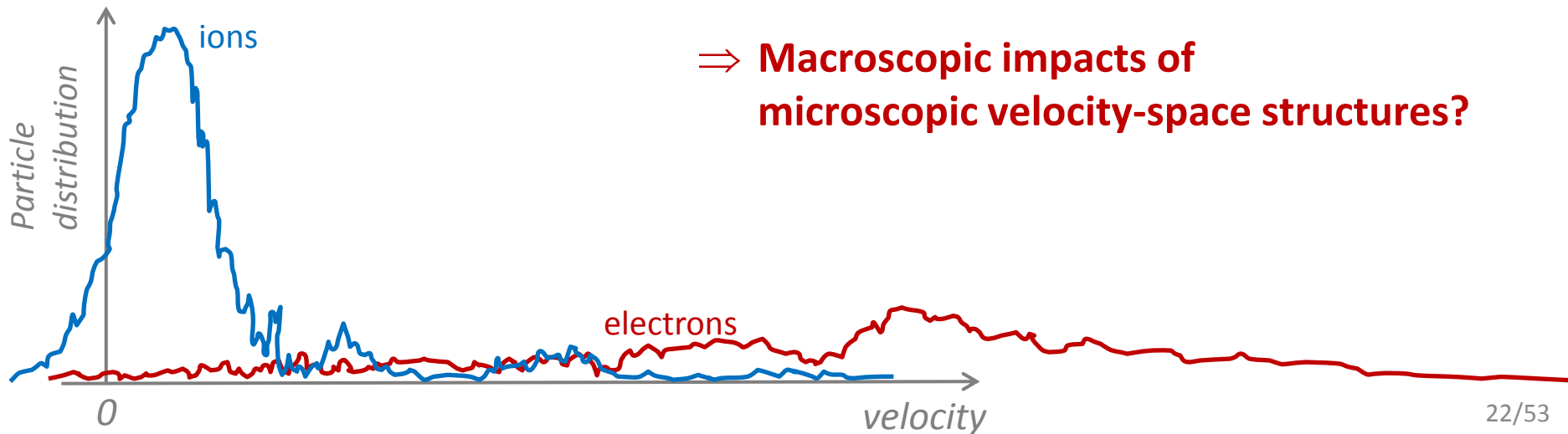
System	Many incompress. fluids in 2D	Quasi-geostrophic fluid	1D collisionless plasma
Distribution	$n(x, y, t)$	$W(x, y, t)$	$f(x, v, t)$
Description space	2D config. space 	2D config. space 	2D phase space 
Continuity equation	$\frac{\partial n}{\partial t} + u_x \frac{\partial n}{\partial x} + u_y \frac{\partial n}{\partial y} = 0$	$\frac{\partial W}{\partial t} + u_x \frac{\partial W}{\partial x} + u_y \frac{\partial W}{\partial y} = 0$	$\frac{\partial f}{\partial t} + v \frac{\partial f}{\partial x} + E \frac{\partial f}{\partial v} = 0$
Self-consistency	Hierarchy of fluid equations + closure	Stream function $W = \nabla^2 \psi$	Poisson $\frac{\partial E}{\partial x} = \sum_s q_s \int f_s dx$

Fluid vs kinetic

1D fluid model: density $n(x, t)$, mean velocity $u(x, t)$, and temperature $T(x, t)$

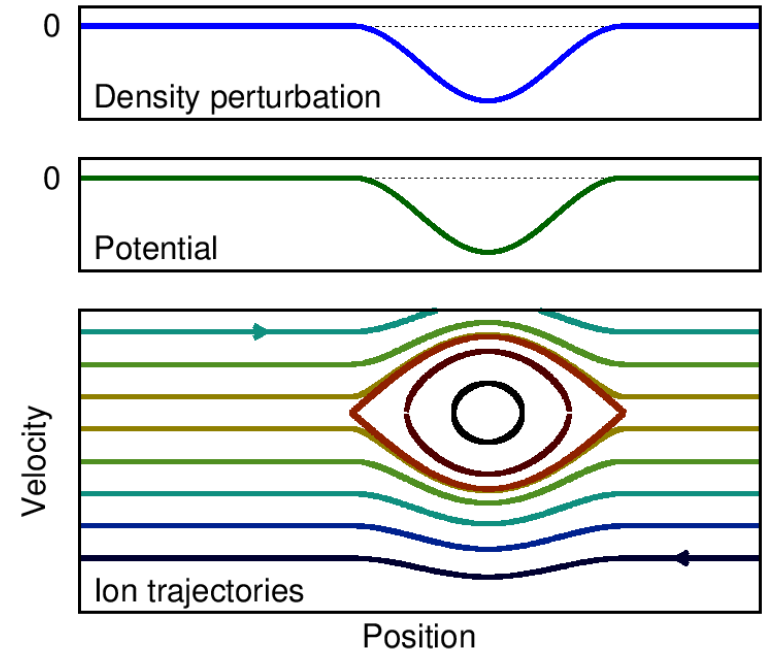
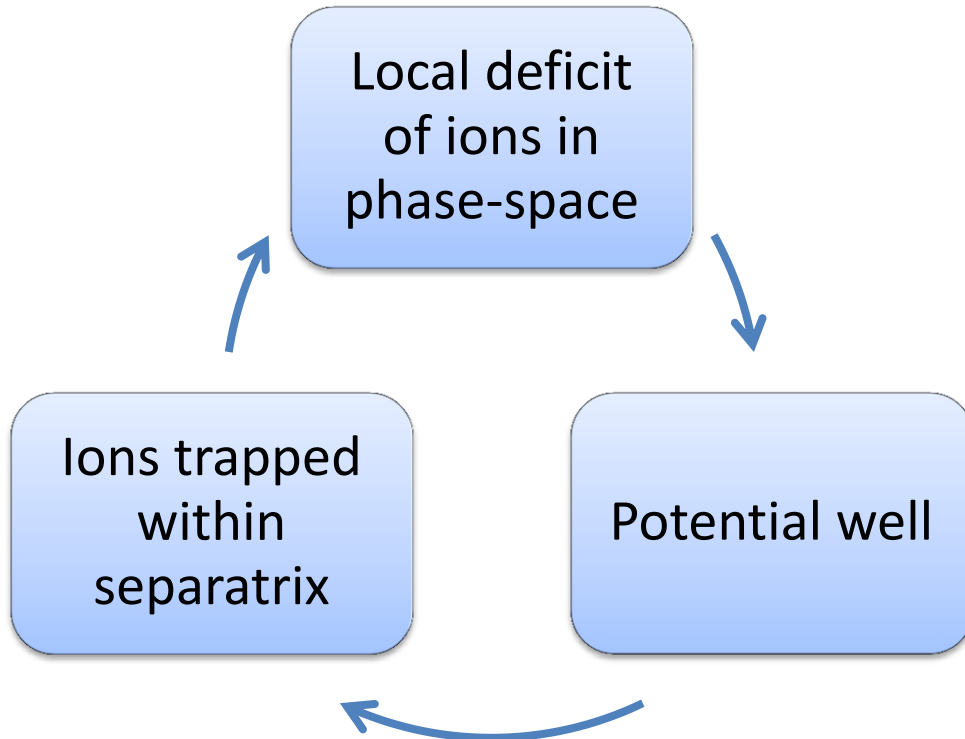


1D kinetic model: fonction de distribution $f(x, v, t)$ \Rightarrow 2D phase-space



Electrostatic trapping yields phase-space vortex

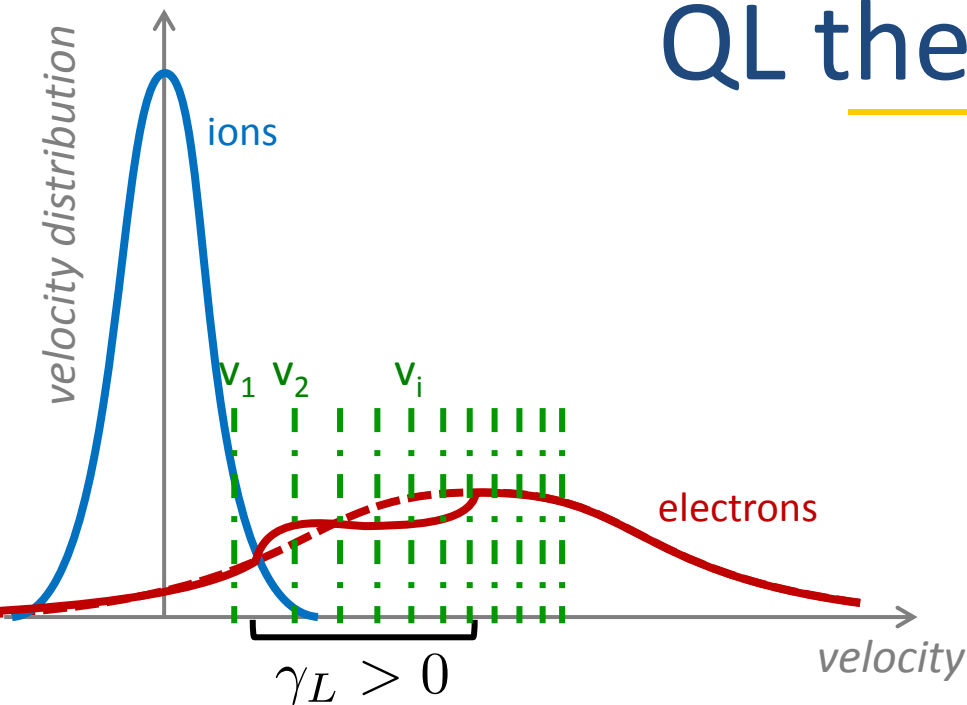
Self-sustaining structure



⇒ BGK mode

Bernstein & Green & Kruskal '57

QL theory: idea



Linear growth rate

$$\gamma_L \sim \left(\frac{1}{m_i} \frac{\partial f_i}{\partial v} + \frac{1}{m_e} \frac{\partial f_e}{\partial v} \right) \Big|_{v=\omega_k/k}$$



Quasilinear theory

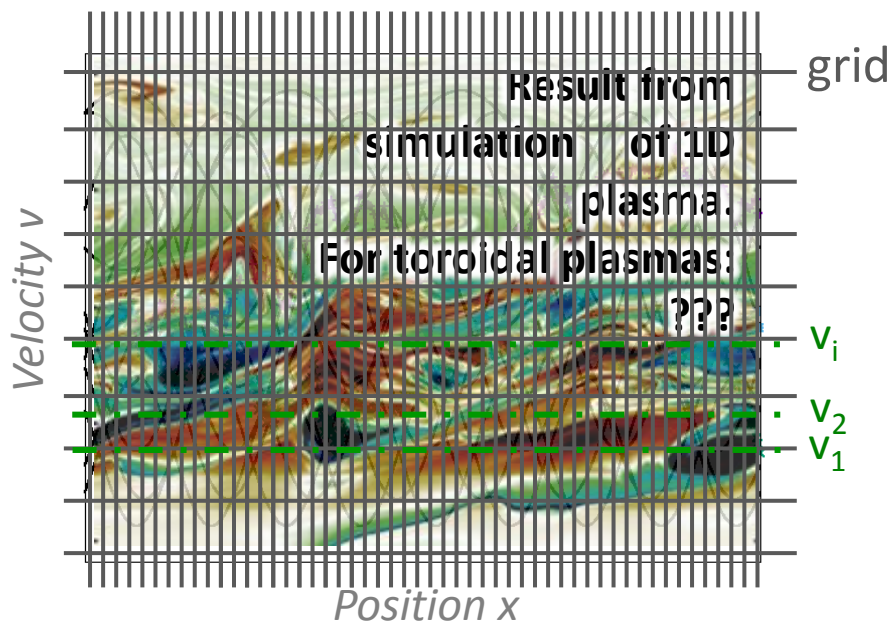
$$\frac{\partial f}{\partial t} = \frac{\partial}{\partial v} \left(D_{\text{QL}} \frac{\partial \tilde{f}}{\partial v} \right)$$

$$D_{\text{QL}} \sim \sum_k |E_k|^2 / k$$

$$\frac{\partial |E_k|}{\partial t} = \gamma_L |E_k|$$



Flattening in the region $\gamma_L > 0$

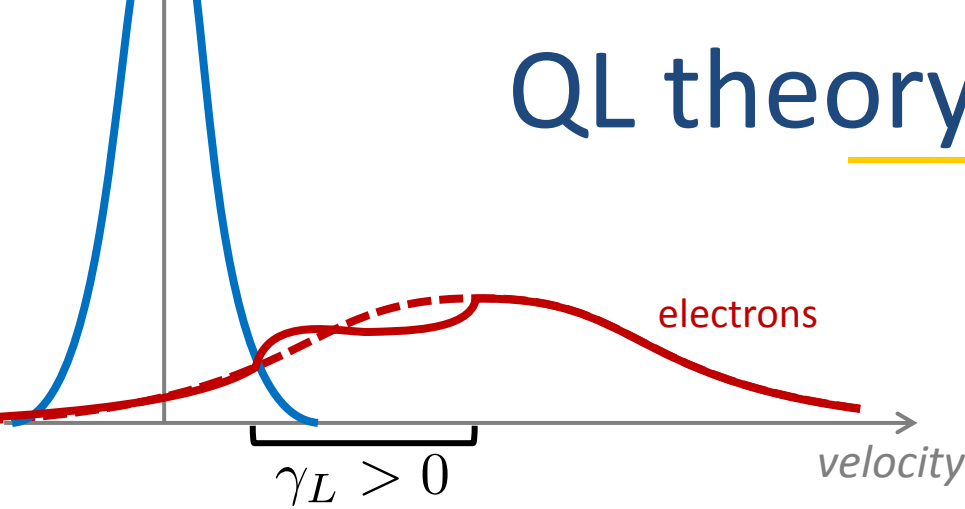


Vedenov, Velikov, Sagdeev '61

Drummond & Pines '62

Sagdeev & Galeev '69

QL theory: limitations

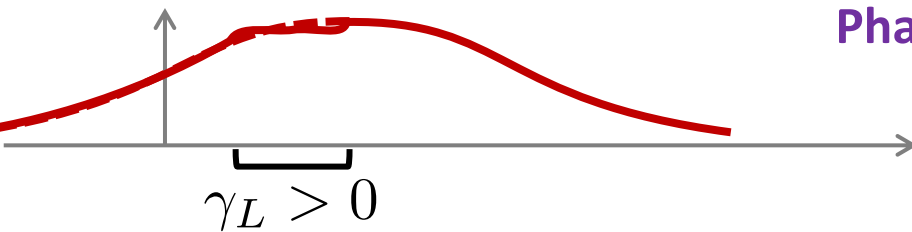


Conventional approach



Flattening in the region $\gamma_L > 0$

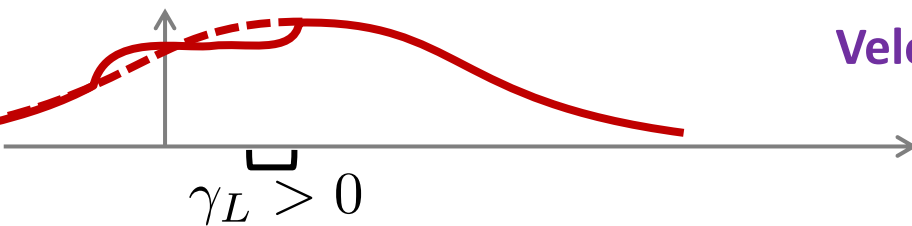
Outside the scope of QL:



Phase-space turbulence

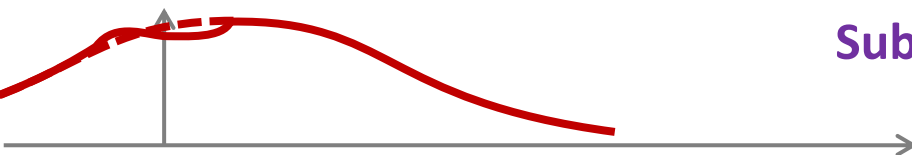
$$\frac{\partial f}{\partial t} = \frac{\partial}{\partial v} \left(F_{\text{PS}} \tilde{f} + D_{\text{PS}} \frac{\partial \tilde{f}}{\partial v} \right)$$

$$D_{\text{PS}} \gg D_{\text{QL}}$$



Velocity-space avalanches

Flattening extends to a region $\gamma_L < 0$



Subcritical instabilities

Heating even though $\gamma_L < 0$ everywhere

Small-scale energy-space structures in GK

Kinetics:

6D Phase-space (3D + 3V)



$$\omega \ll \omega_{\text{gyro}}$$

Gyrokinetics:

4D Phase-space + 1 parameter

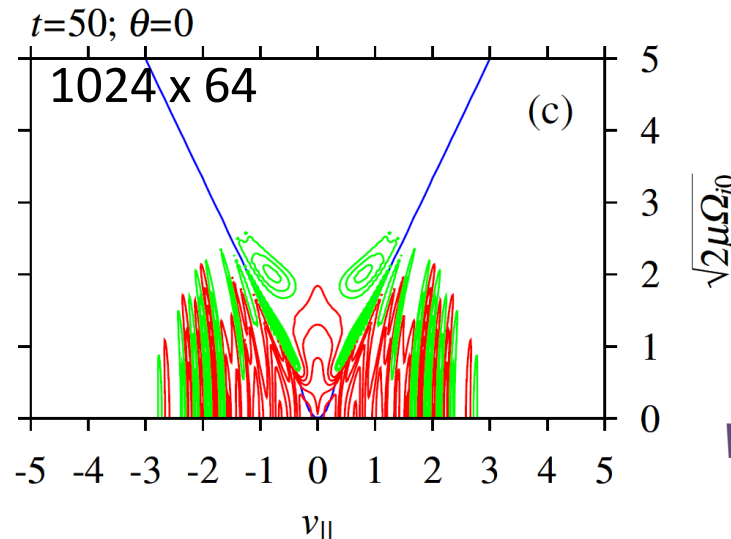
Existence of fine-scale structures in velocity-space

- Not an issue?

Entropy and energy flux constant between $4 \times 4 \times 2$ and $16 \times 16 \times 2$

Candy '06

- Still an issue!



Watanabe '06

⇒ Existing GK simulations may miss effects of small-scale PS structures

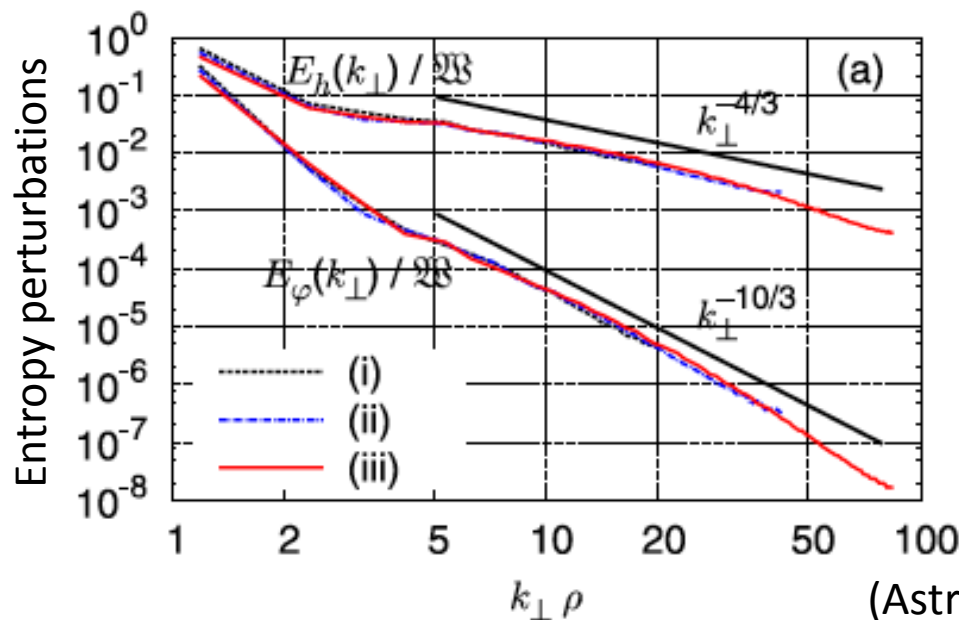
Cascade in velocity-space

Turbulent cascade of entropy in phase-space

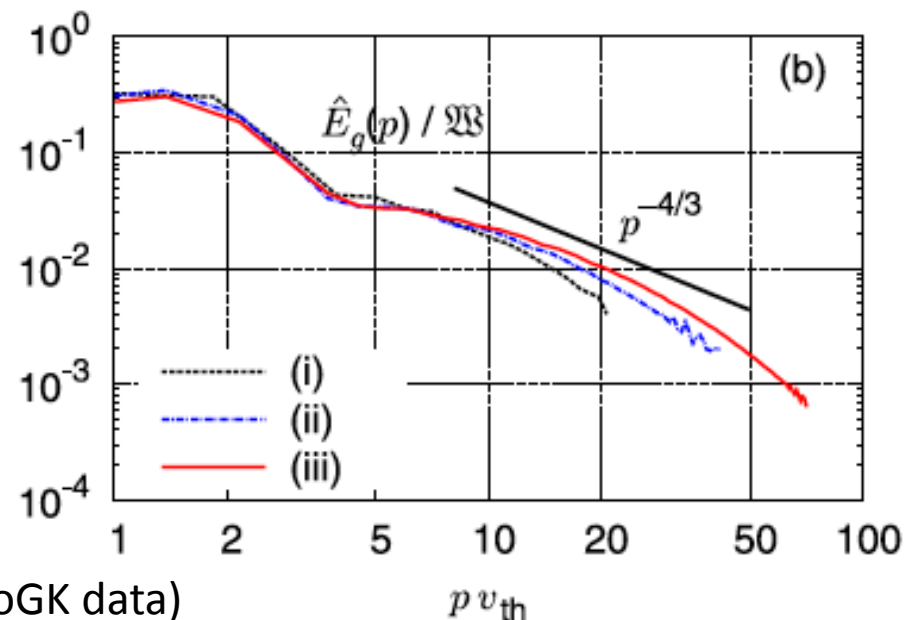
- Driven by nonlinear phase-mixing in strongly turbulent plasmas
- Generation of small-scales in real space and that in velocity space are intertwined.

Schekochihin '08

Tatsuno '09

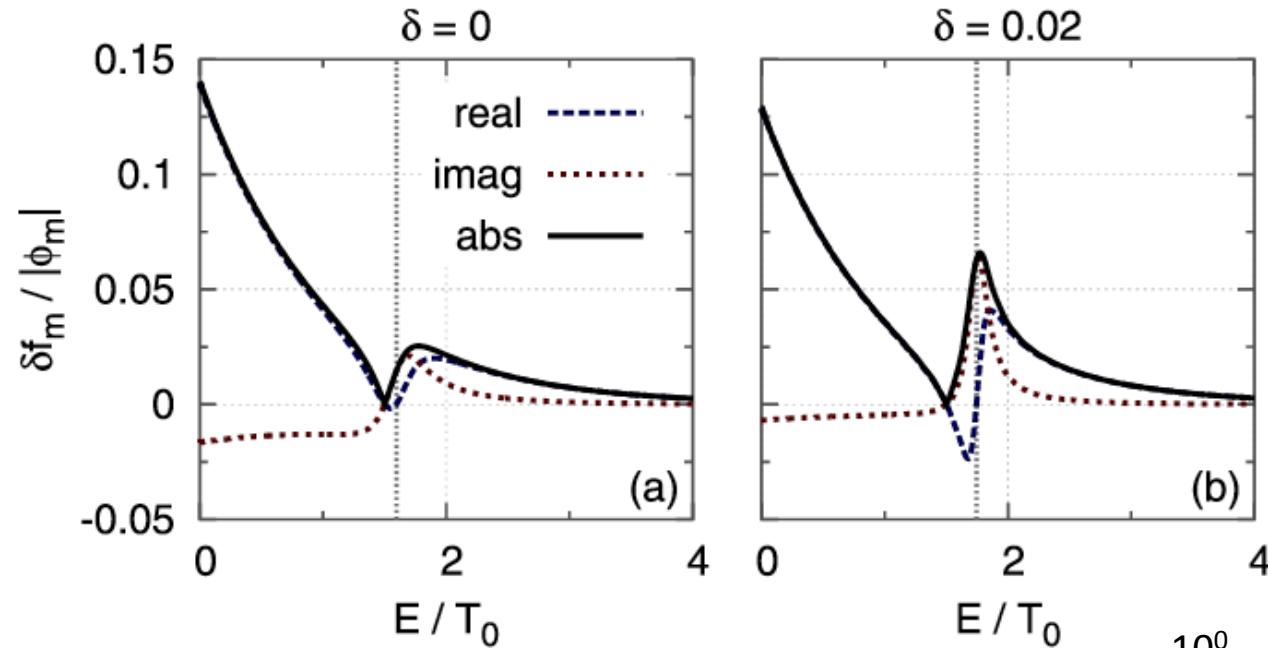


Spectrum in wave number



Spectrum in velocity-space
(Hankel spectrum)

Linear analysis

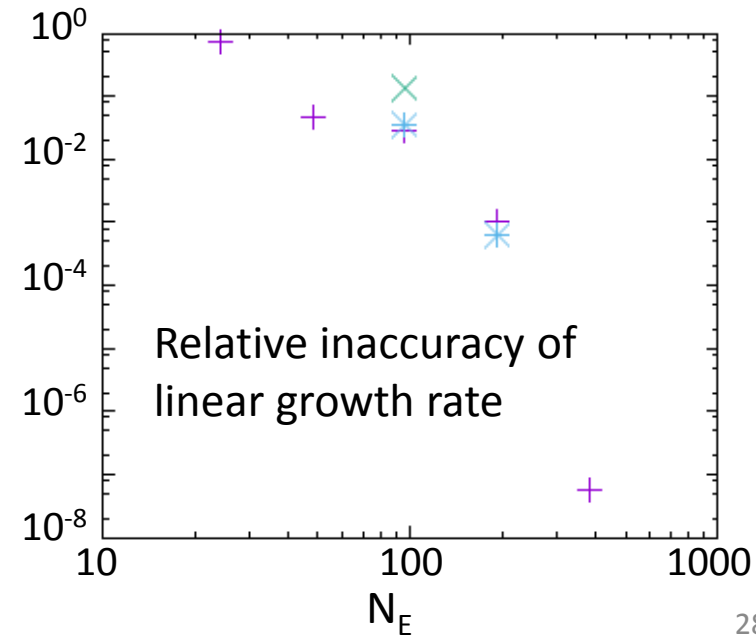


$$\frac{\hat{n}_{e,m}}{n_0} = (1 + \nu m \delta) \frac{e \hat{\phi}_m}{T}$$

$$\delta \approx \epsilon_0^{3/2} \omega_{*e}^- \eta_e / \nu_{ei}$$

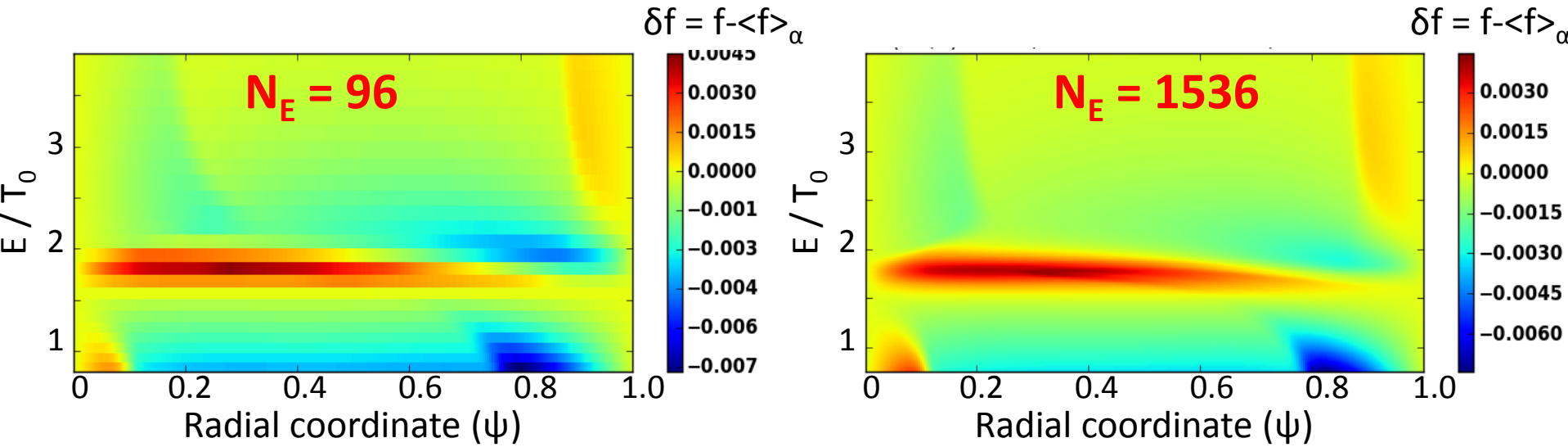
Resonant wave-particle interactions can create a narrow structure in energy-space

- 10% accuracy requires $N_E \approx 200$
- 1% 1000
- Linear growth-rate is not so sensitive (0.1% for $N_E \approx 200$).

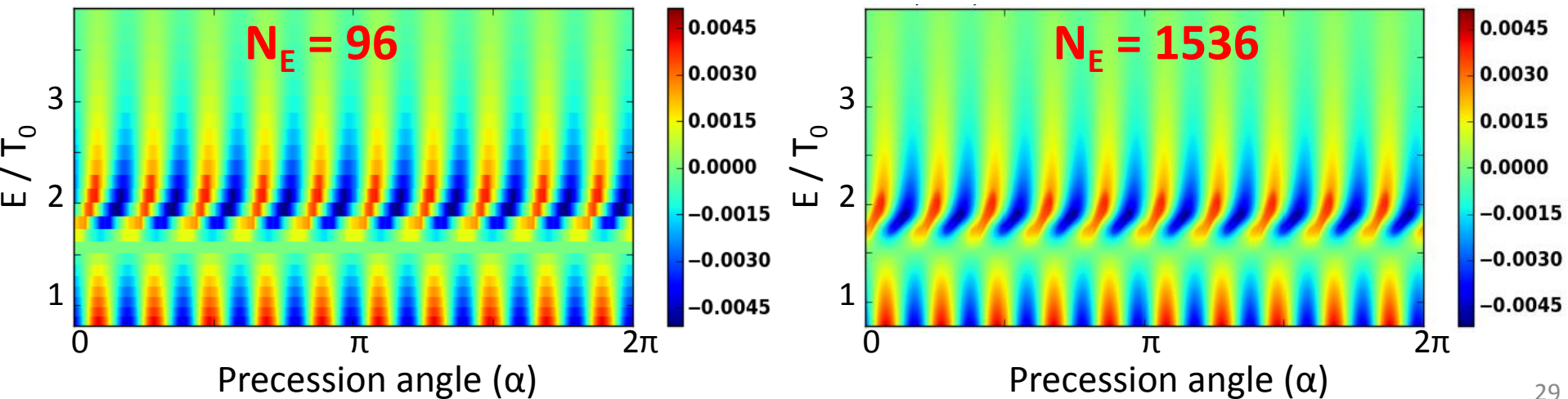


First peak: not-so-fine scales in energy-space

Slice at a given precession angle:

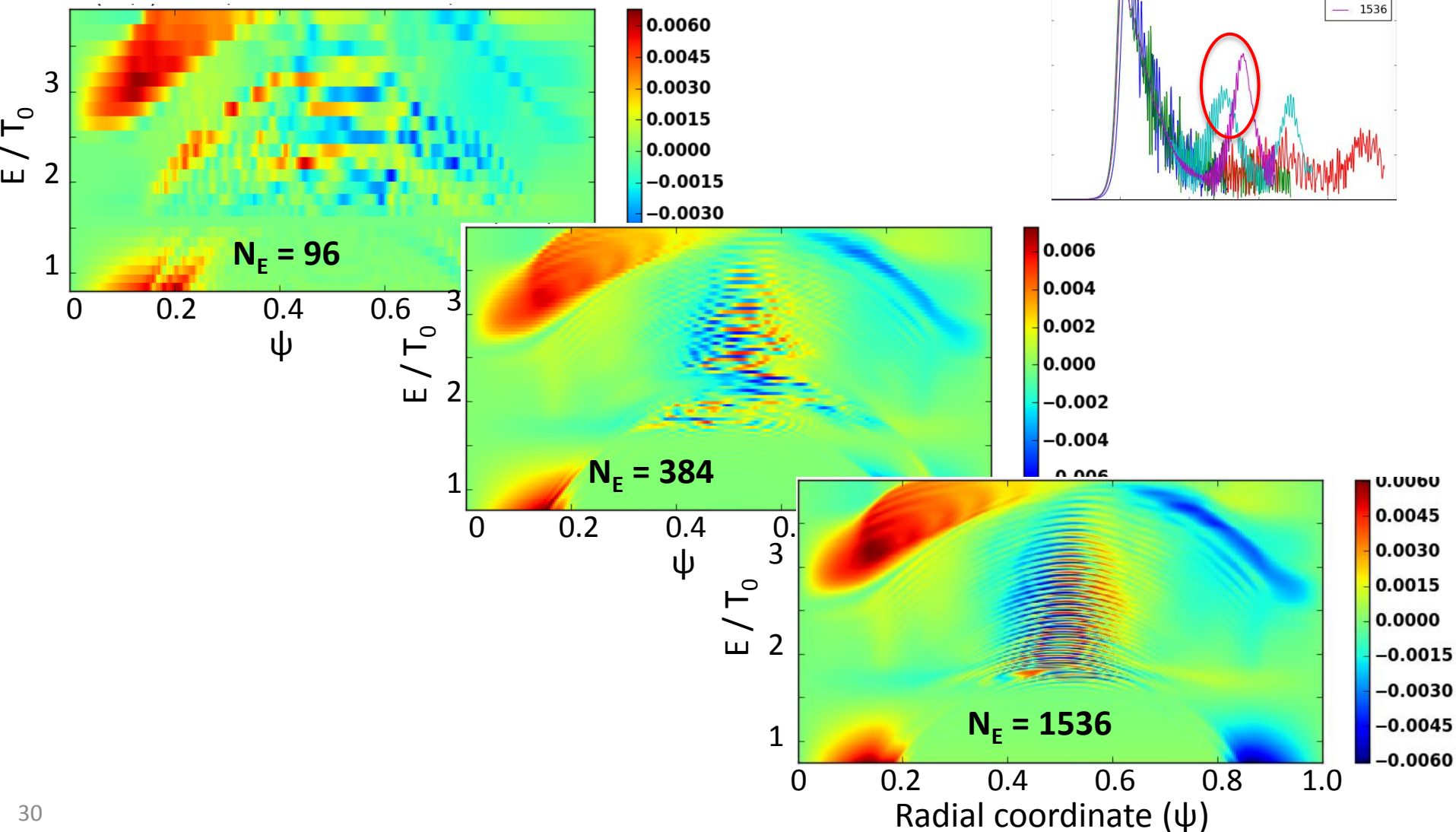


Slice at a given radius:

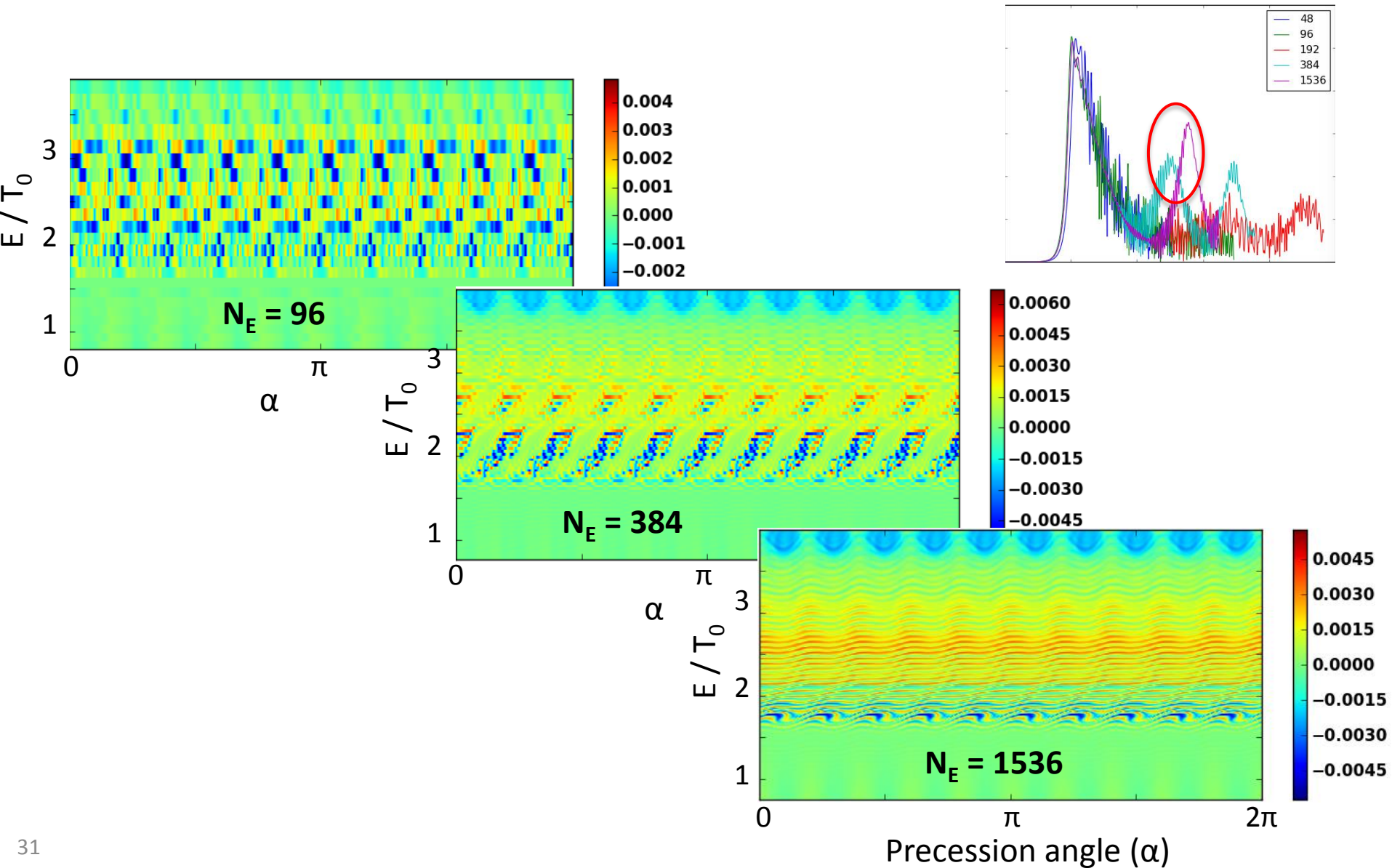


Second peak: slice at a given precession angle

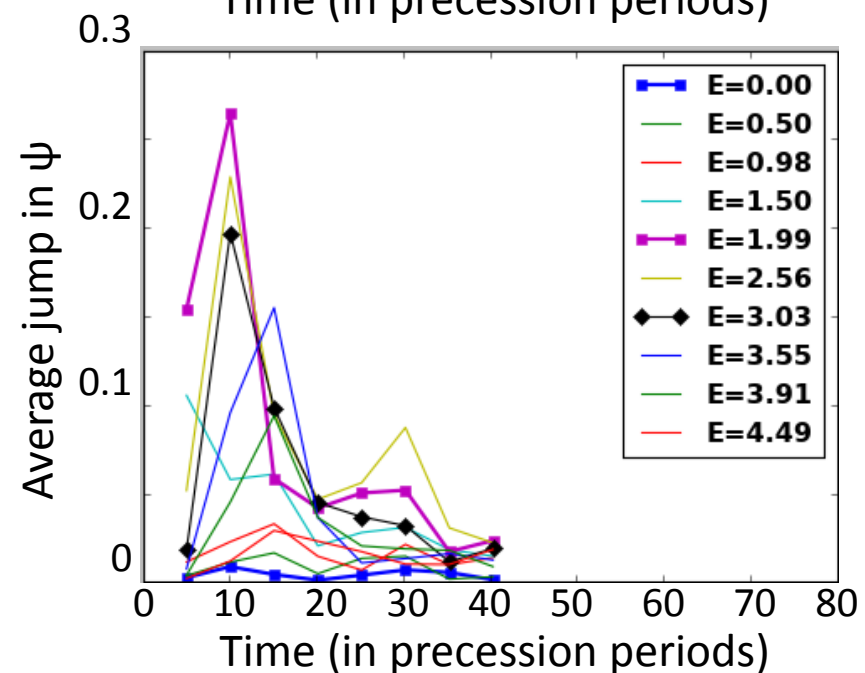
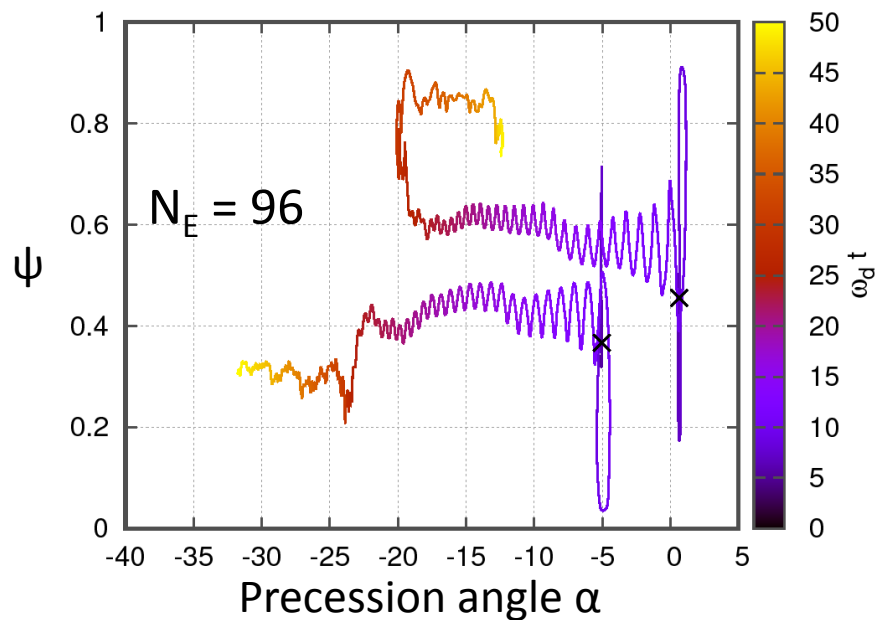
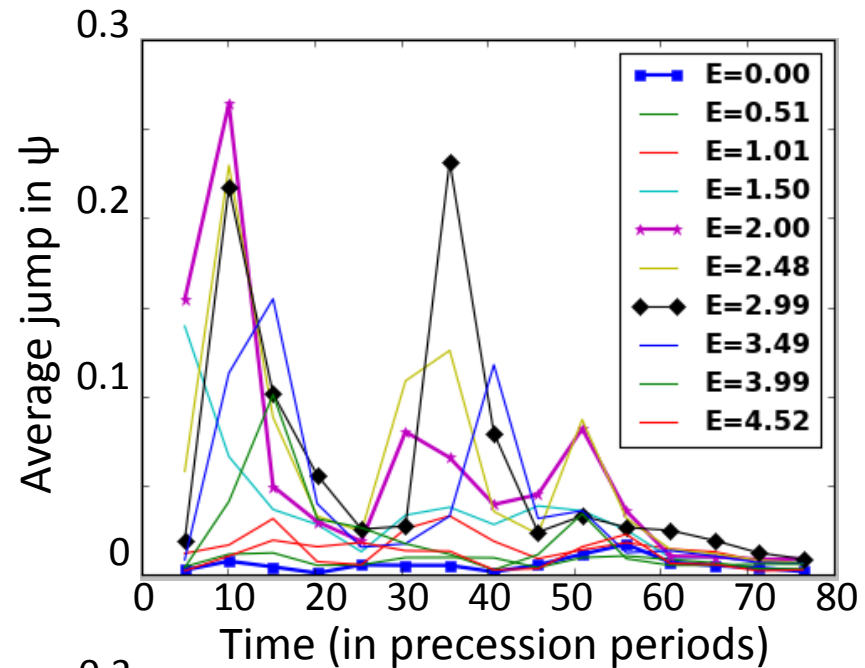
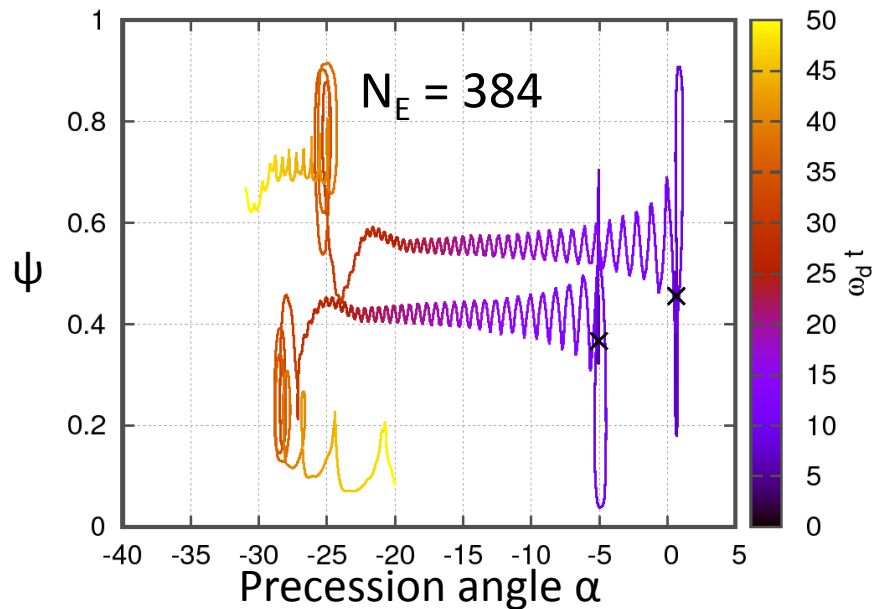
The second peak is associated with the growth of phase-space structures with $\Delta E \ll T_0$



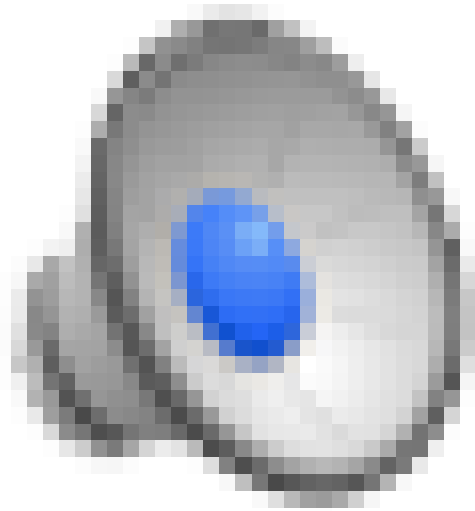
Second peak: slice at a given radius



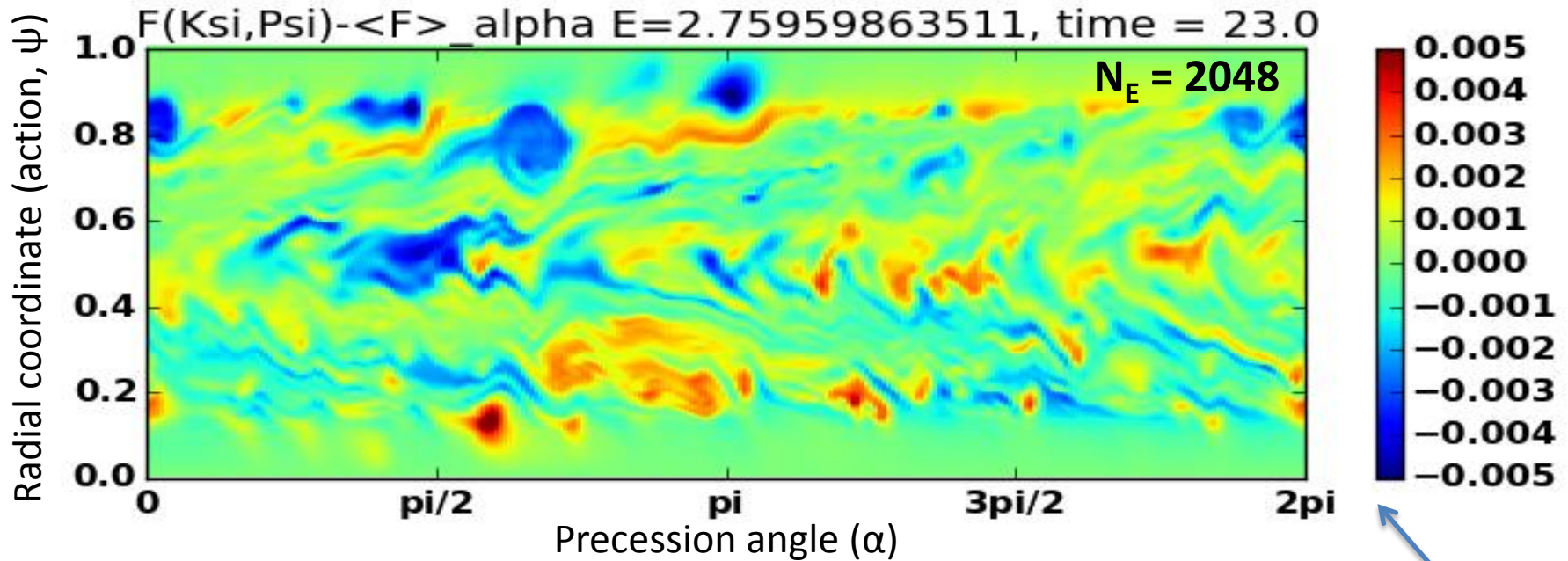
Particle orbits



Radially-dependent precession frequency



These structures survive turbulence



Properties

1. Self-organized from seeds born at the resonant energy.
2. Evolved toward higher energies.
3. Localized in radius, size $\Delta r \sim \rho_{ci} \sim \text{mm}$
4. Localized in energy, with $\Delta E \ll T$
5. Lifetime $1 - 10 \text{ ms} \gg \text{drift period}$

